

Effective Field Theory of Structure Formation

Lecture 2: Lagrangian Approach

z.vlah

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Outline:

1. Dark Matter Clustering in Lagrangian Framework
2. Resummation of the Long Displacements
3. Baryon Acoustic Oscillations
4. Primordial Features in the Power Spectrum
5. Summary

Selected bibliography:

- Large-Scale Galaxy Bias, Desjacques et al., 2018, 1611.09787
- The EFT approach to gravitational dynamics, Porto, 2016, 1601.04914
- Lectures on EFTofLSS, Senatorel, (online notes)
- Modern Cosmology, Dodelson & Schmidt, 2021
- LSS of the Universe and PT, Bernardeau et al., 2002, astro-ph/0112551

Dark Matter Clustering in Lagrangian Framework

Eulerian:



Lagrangian:



Coordinate of a (**t**)racer particle at a given moment in time \mathbf{r}

$$\mathbf{r}(\mathbf{q}, \tau) = \mathbf{q} + \psi(\mathbf{q}, \tau),$$

is given in terms of Lagrangian displacement.

Continuity equation:

$$(1 + \delta(\mathbf{r})) d^3r = d^3q \quad \text{vs.} \quad 1 + \delta(\mathbf{r}) = \int_q \delta^D(\mathbf{r} - \mathbf{q} - \psi(\mathbf{q})),$$

Fourier space

$$(2\pi)^3 \delta^D(\mathbf{k}) + \delta(\mathbf{k}) = \int_q e^{i\mathbf{k} \cdot \mathbf{q}} \exp(i\mathbf{k} \cdot \psi),$$

Lagrangian Perturbation Theory

Density is given by the continuity relation

$$(1 + \delta(\mathbf{x}))d^3x = d^3q.$$

We can combine it with the equation of motion for displacement field:

$$\ddot{\psi} + \mathcal{H}\dot{\psi} = -\nabla\Phi(\mathbf{q} + \psi(\mathbf{q}))$$

and solve perturbatively

$$\psi = \psi^{(1)} + \psi^{(2)} + \psi^{(3)} + \dots$$

Solutions are of the form

$$\psi^{(n)} = \int \frac{d^3p}{(2\pi)^3} L_n(\mathbf{p}_1, \dots, \mathbf{p}_n) \delta_L(\mathbf{p}_1) \dots \delta_L(\mathbf{p}_n).$$

Example of leading kernels solutions

$$L_1(\mathbf{p}_1) = \mathbf{p}_1/p^2$$

$$L_2(\mathbf{p}_1, \mathbf{p}_2) = 3/7(\mathbf{p}_{12})/p_{12}^2 (1 - (\mathbf{p}_1 \cdot \mathbf{p}_2)/(p_1^2 p_2^2))$$

Lagrangian dynamics and EFT

The evolution of ψ is governed by

$$\ddot{\psi} + \mathcal{H}\dot{\psi} = -\nabla\Phi(\mathbf{q} + \psi(\mathbf{q})).$$

Integrating out short modes (using filter $W_R(\mathbf{q}, \mathbf{q}')$) system is splitting into L -long and S -short wavelength modes, e.g.:

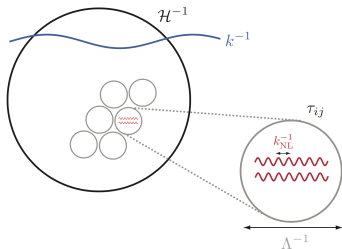
$$\psi_L(\mathbf{q}) = \int_{\mathbf{q}'} W_R(\mathbf{q}, \mathbf{q}')\psi(\mathbf{q}'), \quad \psi_S(\mathbf{q}, \mathbf{q}') = \psi(\mathbf{q}') - \psi_L(\mathbf{q}).$$

This defines δ_L as the long-scale component of the density perturbation corresponding to ψ_L and also Φ_L as the gravitational potential $\nabla^2\Phi_L \sim \delta_L$.
E.o.m. for long displacement:

$$\ddot{\psi}_L + \mathcal{H}\dot{\psi}_L = -\nabla\Phi_L(\mathbf{q} + \psi_L(\mathbf{q})) + \mathbf{a}_S(\mathbf{q}, \psi_L(\mathbf{q})),$$

and $\mathbf{a}_S(\mathbf{q}) = -\nabla\Phi_S(\mathbf{q} + \psi_L(\mathbf{q})) - \frac{1}{2}Q_L^{ij}(\mathbf{q})\nabla\nabla_i\nabla_j\Phi_L(\mathbf{q} + \psi_L(\mathbf{q})) + \dots$,

Similar formalism was also derived in [Porto et al, '14].



Matter power spectrum in Lagrangian EFT

Pair displacement:

$$\Delta = \psi(\mathbf{r}_2) - \psi(\mathbf{r}_1).$$

The correlation function and power spectrum can now be defined through the cumulants of the displacement, e.g.

$$P(k) = \int_q e^{i\mathbf{q}\cdot\mathbf{k}} \left[\langle e^{i\mathbf{k}\cdot\Delta(\mathbf{q})} \rangle - 1 \right].$$

For one loop power spectrum results, keeping linear modes resummed:

$$P(k) = \int_q e^{i\mathbf{k}\cdot\mathbf{q}} \exp \left[-\frac{1}{2} k_i k_j \langle \Delta_i \Delta_j \rangle_c + \frac{i}{6} k_i k_j k_k \langle \Delta_i \Delta_j \Delta_k \rangle_c + \dots \right]$$

Final results equivalent to the Eulerian scheme.

Allows for the insight in the counter term structure and IR resummation schemes (in particular one leads to the scheme in [Senatore+:14]).

Alternative derivation of IR schemes: [Baldauf++:15, Blas++:16].

Lagrangian dynamics and EFT

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For one loop power spectrum results, keeping linear modes resummed:

$$P(k) = \int_q e^{i\mathbf{k}\cdot\mathbf{q} - (1/2)k_i k_j A_{ij}^{\text{lin}}} \left[1 - \frac{1}{2}k_i k_j A_{ij}^{\text{lpt+eft}} + \frac{i}{6}k_i k_j k_k W_{ijk}^{\text{lpt+eft}} + \dots \right]$$

with

$$A_{ij}(\mathbf{q}) = 2 \langle \Psi_i(\mathbf{0}) \Psi_j(\mathbf{0}) \rangle - 2 \langle \Psi_i(\mathbf{q}_1) \Psi_j(\mathbf{q}_2) \rangle$$

Final results equivalent to the Eulerian scheme.

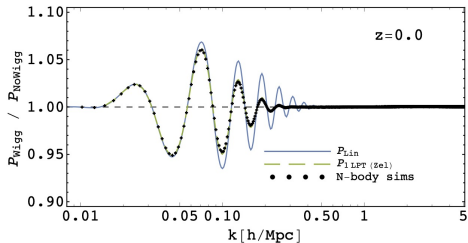
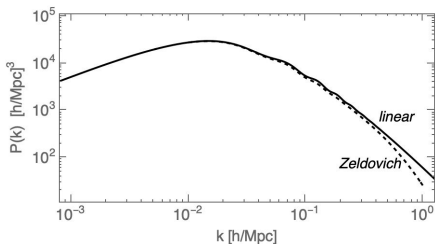
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Zeldovich power spectrum

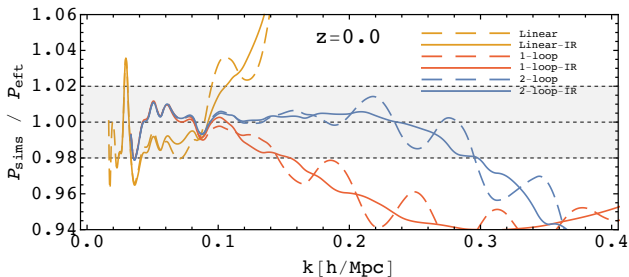
Keeping only linear displacements modes

$$P_{\text{Zel}}(k) = \int_{\mathbf{q}} e^{i\mathbf{k}\cdot\mathbf{q} - (1/2)k_i k_j A_{ij}^{\text{lin}}}$$



Resummation of the Long (IR) Displacements

$$P_{\text{EFT}}(k) = P_0 + P_{1\text{-loop}} + P_{2\text{-loop}} - c_s^2 \frac{k^2}{k_{\text{NL}}^2} (P_{11} + P_{1\text{-loop}}) + \text{c.t.}$$

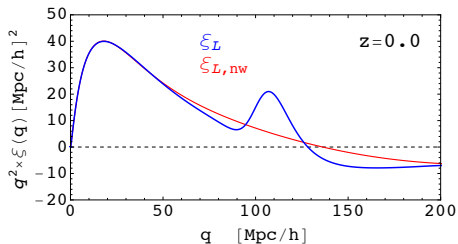
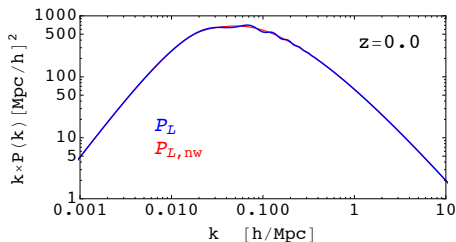


- Well defined/convergent expansion in k/k_{NL} .
- IR divergence and IR safety of equal time correlators (ψ_L).
- IR resummation (Lagrangian approach) - well described BAO

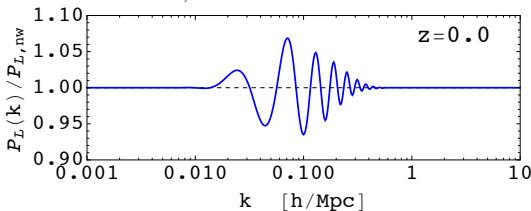
Linear power spectrum, correlation function & BAO

Linear power spectrum P_L : from Boltzmann codes (CAMB, Class).

Can be divided: smooth $P_{L,nw}$ + wiggle part $P_{L,w}$: $P_L = P_{L,nw} + P_{L,w}$



Wiggle power spectrum: $P_{L,w}$.

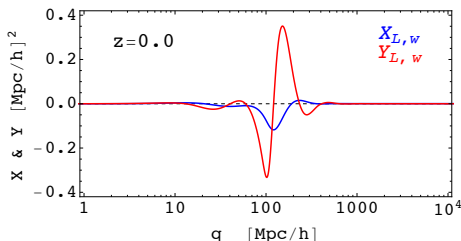
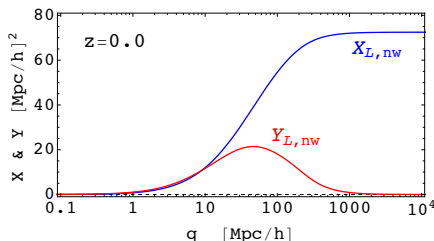


Resummation of IR modes: simple scheme

Separating the wiggle and non-wiggle part

$$A_L^{ij}(\mathbf{q}) = A_{L,nw}^{ij}(\mathbf{q}) + A_{L,w}^{ij}(\mathbf{q});$$

$$P = P_{nw} + \int_q e^{i\mathbf{k}\cdot\mathbf{q} - (1/2)k_i k_j A_{L,nw}^{ij}} \left[-\frac{k_i k_j}{2} \mathcal{A}_{L,w}^{ij} + \dots \right] \simeq P_{nw} + e^{-k^2 \Sigma^2} P_{L,w} + \dots$$

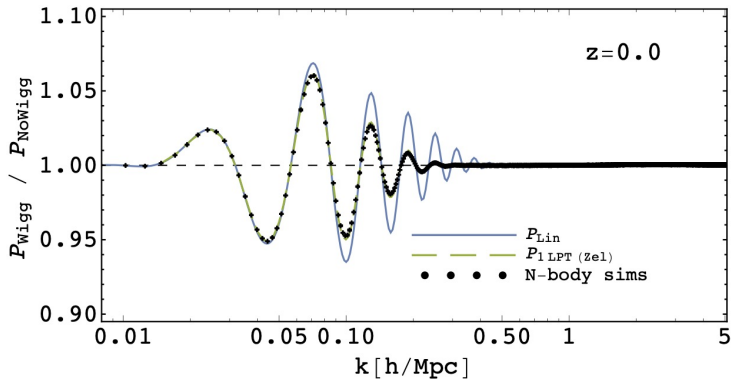


IR-SPT resummation model with $\Sigma^2 = \int \frac{dp}{3\pi^2} (1 - j_0(q_{\max} k)) P_L(p)$:

$$P_{\text{EFT}}(k) = \left[P_0 + P_{1\text{-loop}} - c_s^2 \frac{k^2}{k_{\text{NL}}^2} P_{11} \right. \\ \left. + e^{-k^2 \Sigma^2} \left(P_{w,0} + \Delta P_{w,1\text{-loop}} + (c_s^2/k_{\text{NL}}^2 + \Sigma^2) k^2 P_{w,0} \right) \right]$$

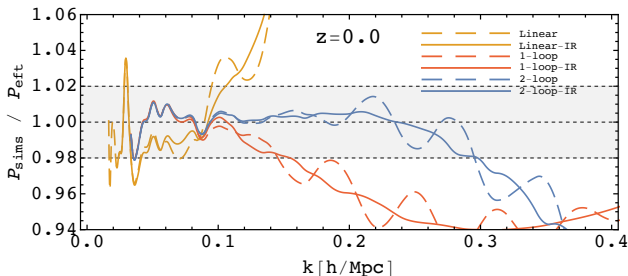
Alternative derivation in: [\[Baldauf:2015\]](#)

BAO wiggles



Power spectrum, loops & BAO

$$P_{\text{EFT}}(k) = P_0 + P_{1\text{-loop}} + P_{2\text{-loop}} - c_s^2 \frac{k^2}{k_{\text{NL}}^2} (P_{11} + P_{1\text{-loop}}) + \text{c.t.}$$



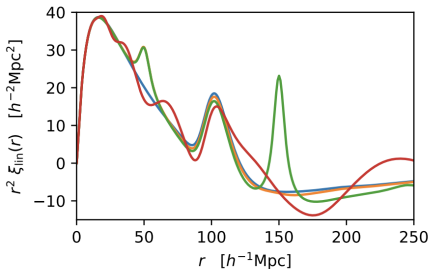
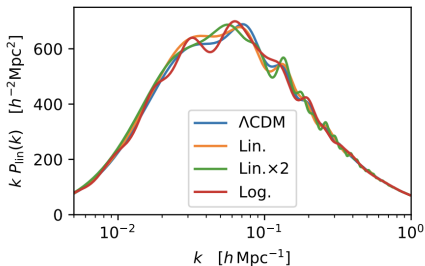
- Well defined/convergent expansion in k/k_{NL} .
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Primordial Features in the Power Spectrum

Phenomenological models of primordial features:

$$\text{(linear)} : P_L = P_{\Lambda\text{CDM}} \left(1 + A \sin(\omega k) \exp(-(kr_d)^2) \right)$$

$$\text{(logarithmic)} : P_L = P_{\Lambda\text{CDM}} \left(1 + A \sin(\omega \ln(k/k^*)) \exp(-(kr_d)^2) \right)$$



Primordial Features in the Power Spectrum

Power spectrum after resummation of the IR modes.

