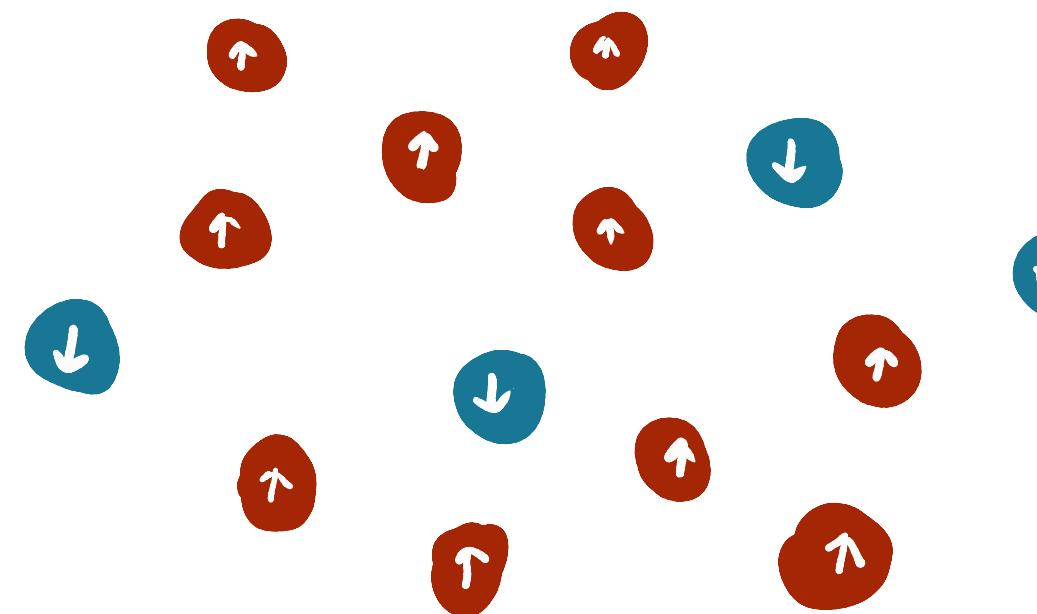


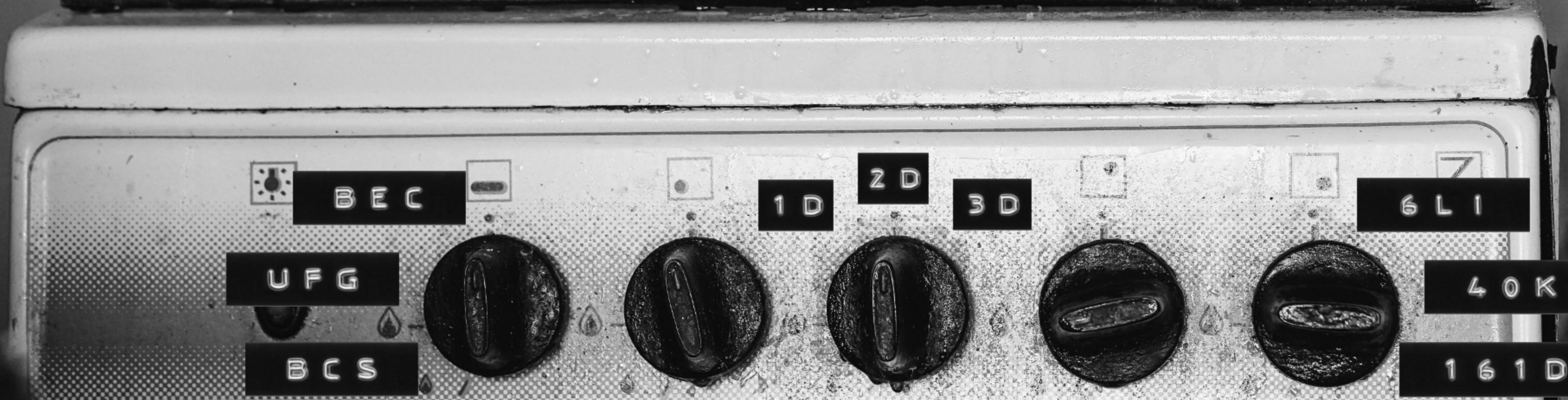
unitary fermions with finite spin-asymmetry

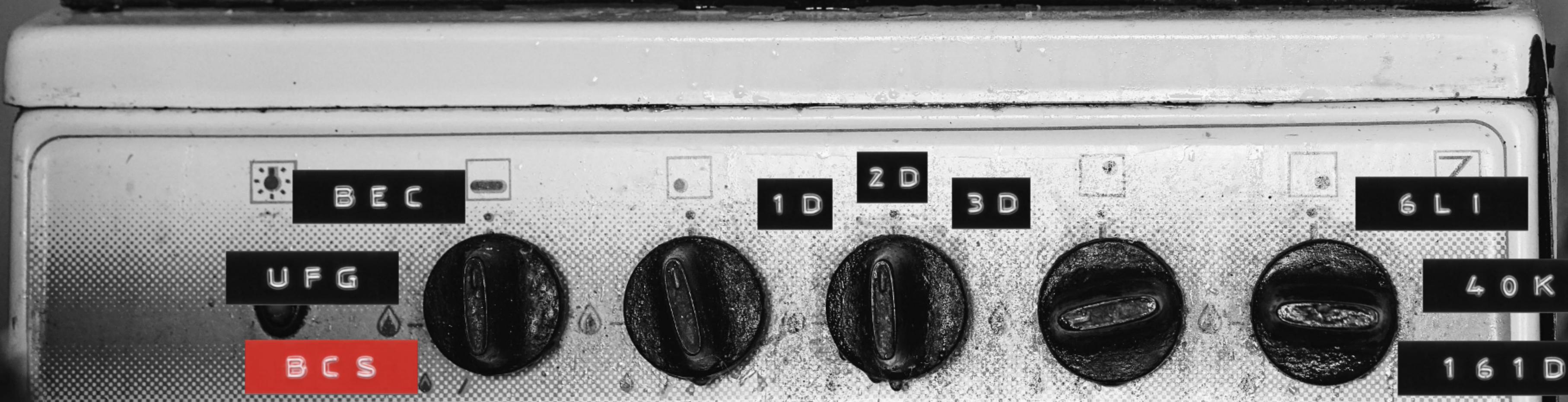
[Berger, Loheac, LR, Ehmann, Braun, Drut *arXiv:1907.10183*
[LR, Loheac, Drut, Braun *Phys. Rev. Lett.* 121, 173001, 2018]

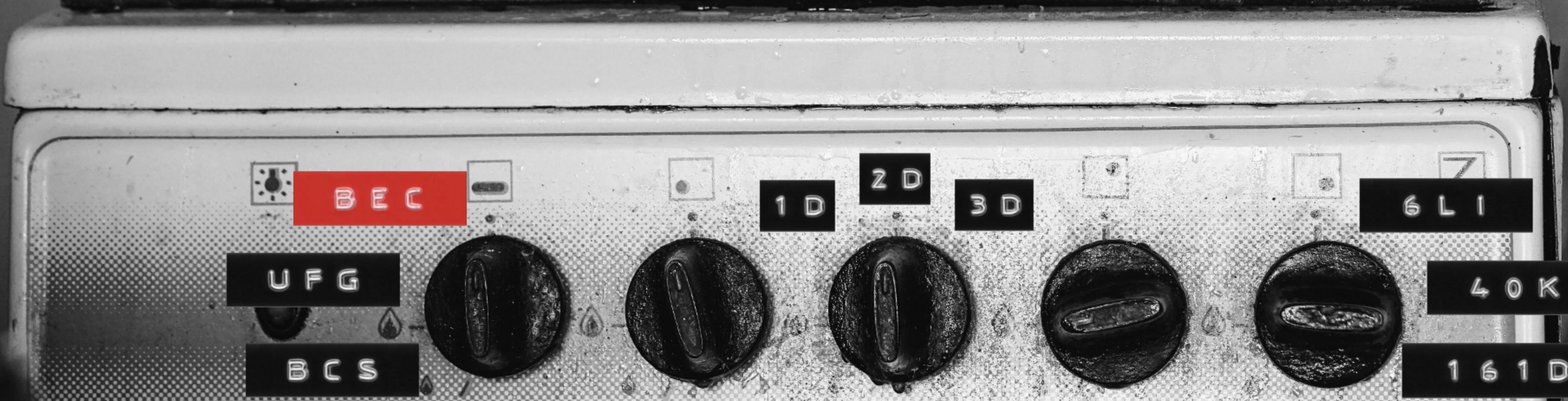


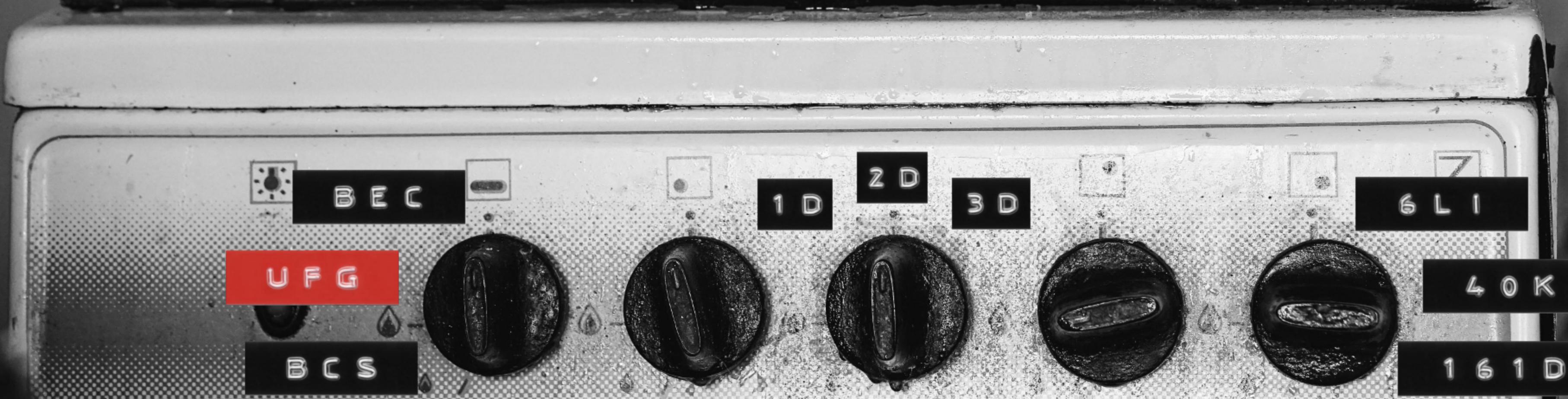
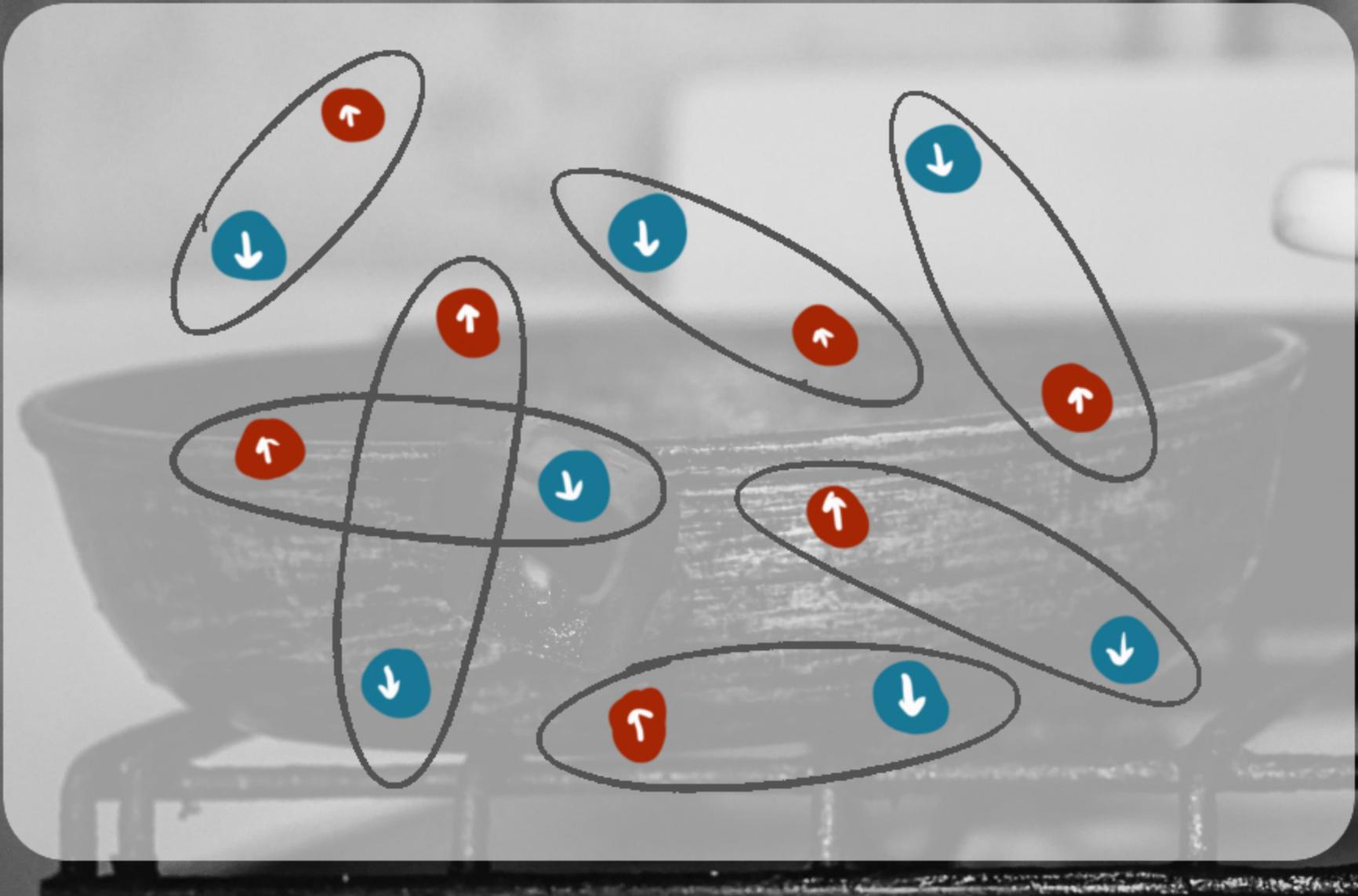
Lukas Rammelmüller, TU Darmstadt

Quantum Systems in Extreme Conditions - Heidelberg, September 24, 2019









the unitary Fermi gas (UFG)

[reviews: Zwerger '12; Mukaiyama,Ueda '13]

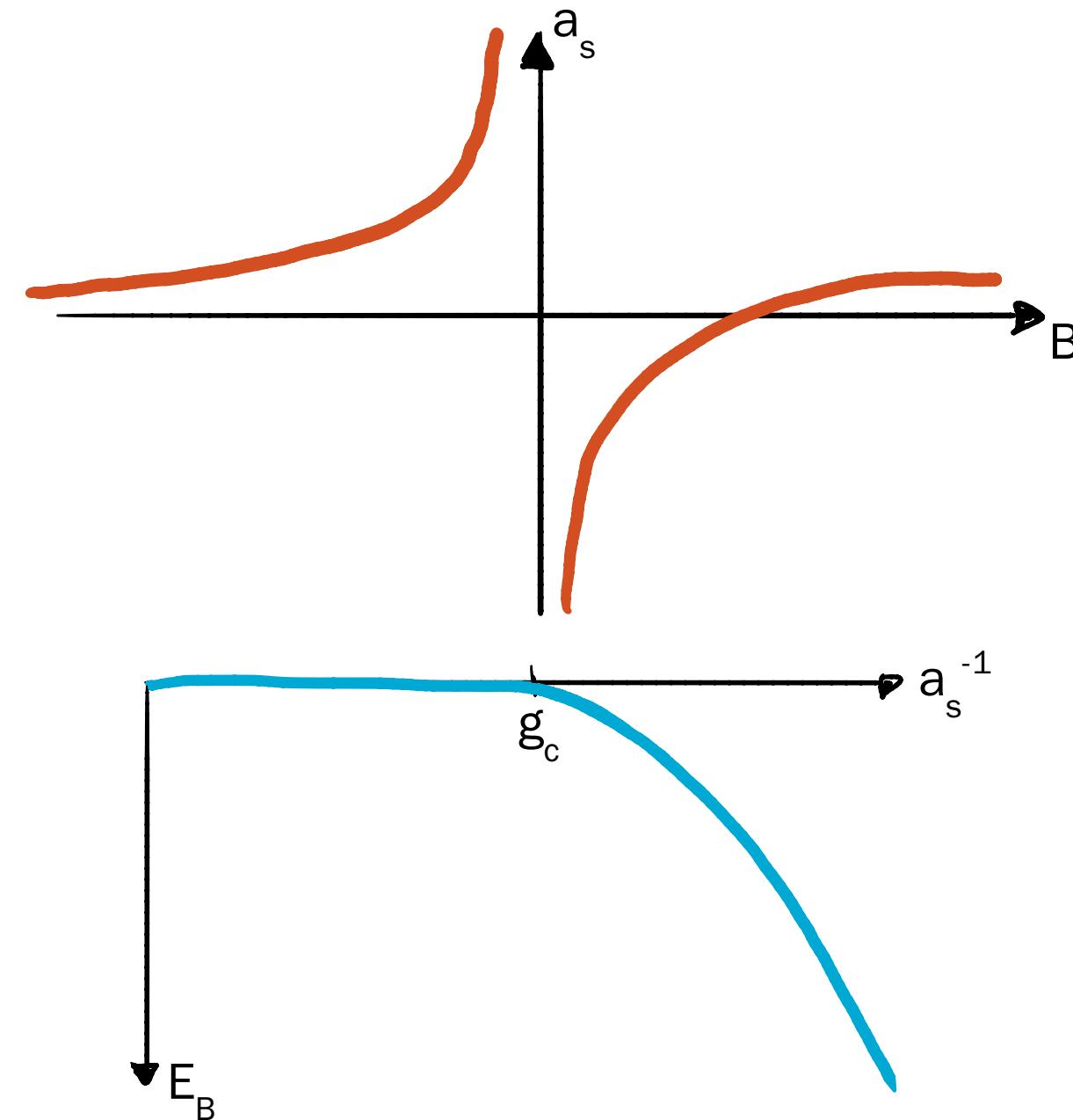
$$a_s \gg n^{-1/3} \gg r_0$$

- ▶ density & temperature are the **only dimensionful** scales in the system
- ▶ **universal** scaling functions:

$$E = E_{FG} f_E(n, T)$$

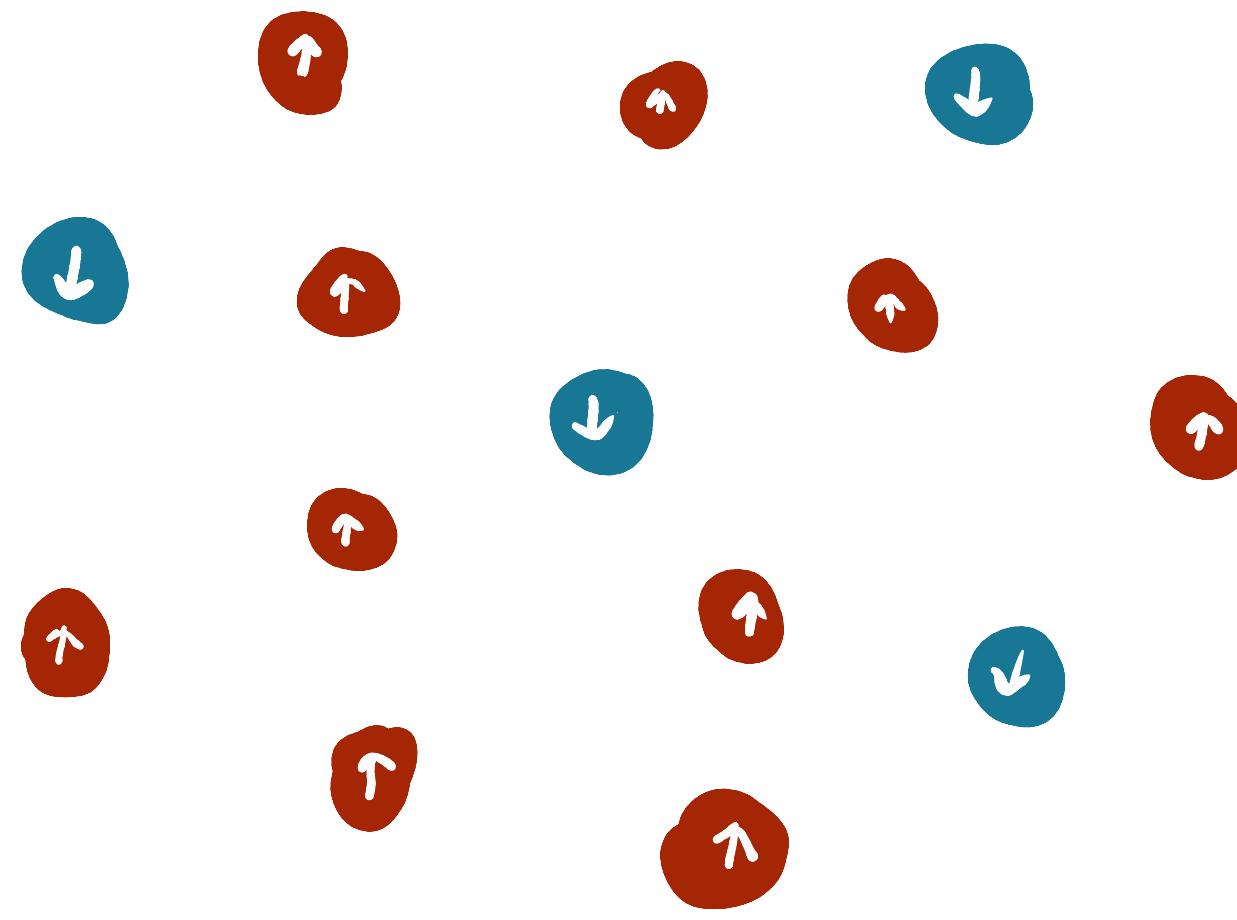
numerous experiments:

- first realizations of unitary fermions [Regal,Greiner,Jin '04; Zwierlein et al. '04; Kinast et al. '04]
- universal behavior & thermodynamics [Thomas,Kinast,Turpalov '05; Horikoshi et al. '10]
- temperature vs. polarization phase-diagram [Shin,Schunck,Schirotzek,Ketterle '08]
- measurement of equation of state [Nascimbène et al. '10; van Houcke et al. '12]
- superfluid transition [Ku,Sommer,Cheuck,Zwierlein '12]
- temperature dependence of Tan's contact [Carcy et al. '19; Mukherjee et al. '19]
- and many more...

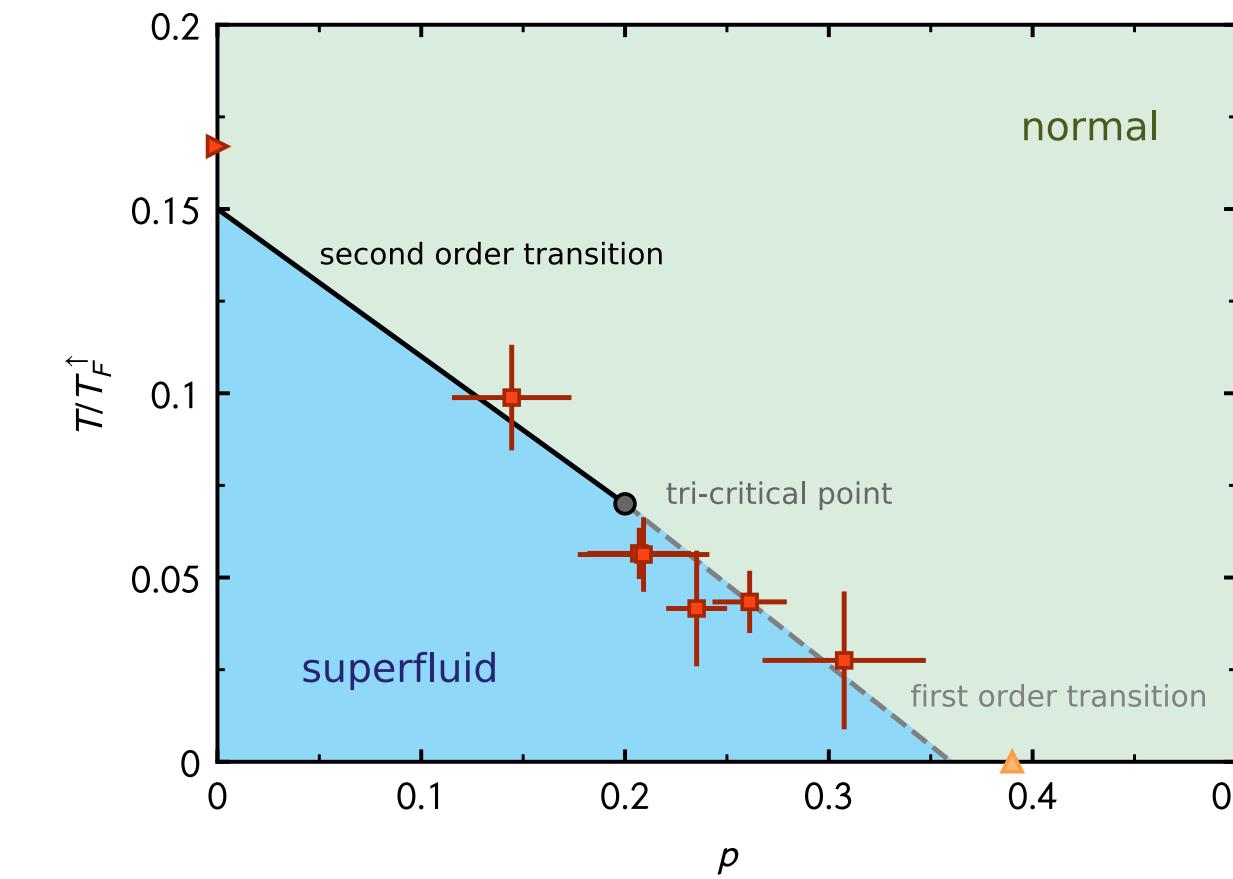


finite spin polarization

[reviews: Chevy,Mora '10; Gubbels,Stoof '13]



[experimental PD: Shin,Schunck,Schirotzek,Ketterle '08]



Many questions remain:
critical polarization? phase diagram? pairing?

agenda

part I

quick intro to **stochastic quantization & CL**
(what is it & how can it help us for fermions?)

part II

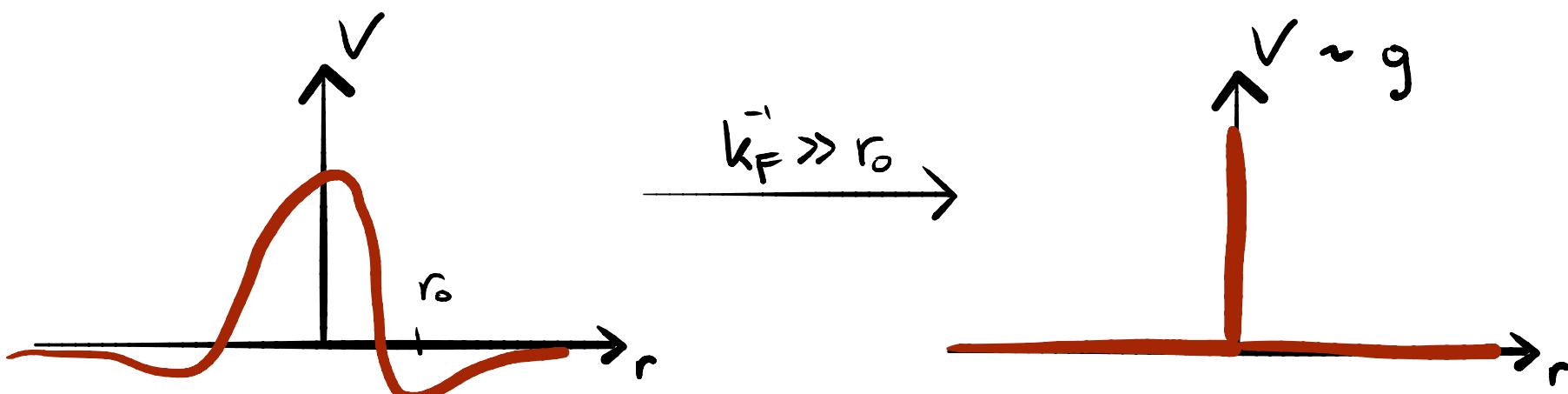
unitary fermions with finite polarization
(equations of state & thermodynamic response)

fermions with contact interaction

$$\hat{H} = - \sum_{s=\uparrow,\downarrow} \int d^d x \hat{\psi}_s^\dagger(\vec{x}) \left(\frac{\hbar^2 \vec{\nabla}^2}{2m_s} \right) \hat{\psi}_s(\vec{x}) + g \int d^d x \hat{\psi}_\uparrow^\dagger(\vec{x}) \hat{\psi}_\uparrow(\vec{x}) \hat{\psi}_\downarrow^\dagger(\vec{x}) \hat{\psi}_\downarrow(\vec{x})$$

kinetic part

interaction part



what do we need to compute?

[lattice methods: Lee '09; Drut,Nicholson '13; Zhang '13]

$$\mathcal{Z} = \text{Tr}[\text{e}^{-\beta \hat{H}}] = \text{Tr}[\text{e}^{-\beta(\hat{T} + \hat{V})}]$$

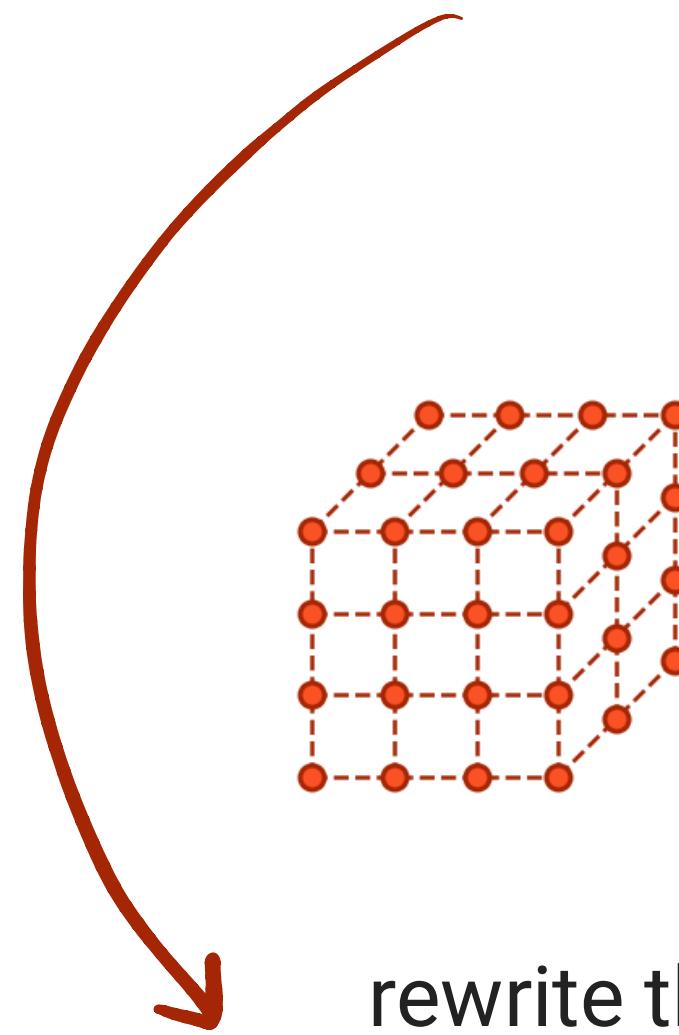
$$\langle \mathcal{O} \rangle = \frac{1}{\mathcal{Z}} \text{Tr}[\hat{\mathcal{O}} \text{e}^{-\beta \hat{H}}]$$

what do we need to compute?

[lattice methods: Lee '09; Drut,Nicholson '13; Zhang '13]

$$\mathcal{Z} = \text{Tr}[e^{-\beta \hat{H}}] = \text{Tr}[e^{-\beta(\hat{T} + \hat{V})}]$$

$$\langle \mathcal{O} \rangle = \frac{1}{\mathcal{Z}} \text{Tr}[\hat{\mathcal{O}} e^{-\beta \hat{H}}]$$



+ Trotter decomposition

+ Hubbard-Stratonovich
transformation

rewrite the problem as a **path-integral**:

$$\mathcal{Z} = \int \mathcal{D}\phi \det M_\phi^\uparrow \det M_\phi^\downarrow \equiv \int \mathcal{D}\phi e^{-S[\phi]}$$

(hard problem = \sum easy problems)

stochastic quantization

[Parisi, Wu '81; Damgaard, Hüffel '87]

$$\mathcal{Z} = \int \mathcal{D}\phi e^{-S[\phi]}$$

(partition function)

$$\langle \hat{\mathcal{O}} \rangle = \int \mathcal{D}\phi \mathcal{O}[\phi] P[\phi]$$

(expectation values)

$$P[\phi] = \frac{1}{\mathcal{Z}} e^{-S[\phi]}$$

(probability measure)

key idea:

probability measure of a **d-dimensional Euclidean path integral**
as equilibrium distribution of a **d+1-dimensional random process**

stochastic quantization

[Parisi, Wu '81; Damgaard, Hüffel '87]

random walk governed by

Langevin equation (Brownian motion):

$$\frac{\partial \phi}{\partial t_L} = - \frac{\delta S[\phi]}{\delta \phi} + \eta$$

stochastic quantization

[Parisi, Wu '81; Damgaard, Hüffel '87]

random walk governed by

Langevin equation (Brownian motion):

$$\frac{\partial \phi}{\partial t_L} = - \frac{\delta S[\phi]}{\delta \phi} + \eta$$



**fictitious Langevin time
(not physical)**

stochastic quantization

[Parisi, Wu '81; Damgaard, Hüffel '87]

random walk governed by

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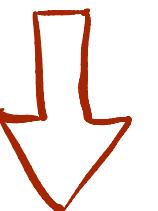
fictitious Langevin time
(not physical)

noise term
 $\langle \eta \rangle = 0$
 $\langle \eta_t \eta_{t'} \rangle = 2\delta(t - t')$

the Langevin method

[Parisi, Wu '81; Damgaard, Hüffel '87]

$$\frac{\partial \phi}{\partial t_L} = -\frac{\delta S[\phi]}{\delta \phi} + \eta$$

 **discretization**

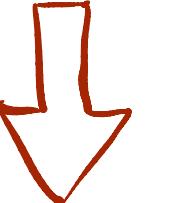
$$\phi^{(n+1)} = \phi^{(n)} - \frac{\delta S[\phi]}{\delta \phi} \Big|_{\phi^{(n)}} \Delta t_L + \sqrt{2\Delta t_L} \tilde{\eta}$$

(Markov chain)

the Langevin method

[Parisi, Wu '81; Damgaard, Hüffel '87]

$$\frac{\partial \phi}{\partial t_L} = -\frac{\delta S[\phi]}{\delta \phi} + \eta$$

 **discretization**

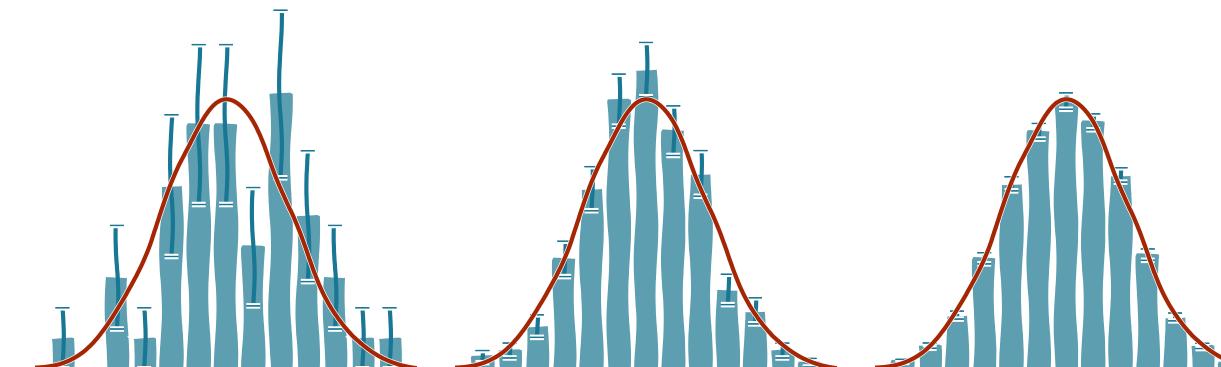
$$\phi^{(n+1)} = \phi^{(n)} - \frac{\delta S[\phi]}{\delta \phi} \Big|_{\phi^{(n)}} \Delta t_L + \sqrt{2\Delta t_L} \tilde{\eta}$$

(Markov chain)

statistical evaluation

$$\langle \hat{\mathcal{O}} \rangle = \frac{1}{N} \sum_{i=1}^N \mathcal{O}[\phi_i]$$

$$\sigma \propto \left(\sqrt{\# \text{ of (uncorrelated) samples}} \right)^{-1}$$



how can stochastic quantization
help us to study fermions?

$$\mathcal{Z} = \int \mathcal{D}\phi \det M_\phi^\uparrow \det M_\phi^\downarrow = \int \mathcal{D}\phi e^{-S[\phi]}$$

how can stochastic quantization help us to study fermions?

$$\mathcal{Z} = \int \mathcal{D}\phi \det M_{\phi}^{\uparrow} \det M_{\phi}^{\downarrow} = \int \mathcal{D}\phi e^{-S[\phi]}$$

probability measure **not positive (semi-)definite**
if any of these conditions applies:

$$\begin{aligned}N_{\uparrow} &\neq N_{\downarrow} \\ \mu_{\uparrow} &\neq \mu_{\downarrow} \\ m_{\uparrow} &\neq m_{\downarrow} \\ g &> 0\end{aligned}$$

the path integral & complex actions

$$\mathcal{Z} = \int \mathcal{D}\phi \det M_\phi^\uparrow \det M_\phi^\downarrow = \int \mathcal{D}\phi e^{-S[\phi]}$$

discrete Langevin equation:

$$\phi^{(n+1)} = \phi^{(n)} + \Delta\phi^{(n)}$$

$$\Delta\phi_R^{(n)} = -\text{Re} \left[\frac{\delta S[\phi]}{\delta \phi} \right]_{\phi^{(n)}} \Delta t_L + \sqrt{2\Delta t_L} \eta$$

$$\Delta\phi_I^{(n)} = -\text{Im} \left[\frac{\delta S[\phi]}{\delta \phi} \right]_{\phi^{(n)}} \Delta t_L$$

complex action → complex Langevin equation

[Parisi '83; Klauder '84; Koonin,Adami '01; Berges,Stamatescu '05; Aarts '08; LR,Porter,Drut,Braun '17; LR,Loheac,Drut,Braun '18; Berger et al. '19]

complex probabilities

$$\int \mathcal{D}\phi P[\phi] O[\phi] \stackrel{?}{=} \int \mathcal{D}\phi_R \mathcal{D}\phi_I P[\phi_R, \phi_I] O[\phi_R + i\phi_I]$$

guaranteed convergence if PDF decays
fast enough and $S[\phi]$ is holomorphic

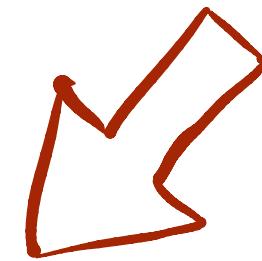
[Aarts,Seiler,Stamatescu '10; Aarts,James,Seiler,Stamatescu '11]

complex probabilities & possible issues

$$\int \mathcal{D}\phi P[\phi] O[\phi] \stackrel{?}{=} \int \mathcal{D}\phi_R \mathcal{D}\phi_I P[\phi_R, \phi_I] O[\phi_R + i\phi_I]$$

guaranteed convergence if PDF decays
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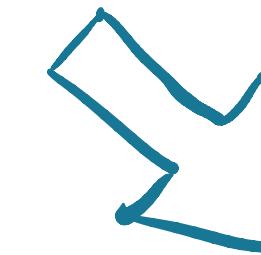
[Aarts,Seiler,Stamatescu '10; Aarts,James,Seiler,Stamatescu '11]



non-analyticities in the action

- zeros in measure ($\det M = 0$)
- could lead to ergodicity issues (bottlenecks)

[Aarts,Seiler,Sexty,Stamatescu '17]



non-vanishing boundary terms

- convergence to wrong limits possible
- behavior must be monitored

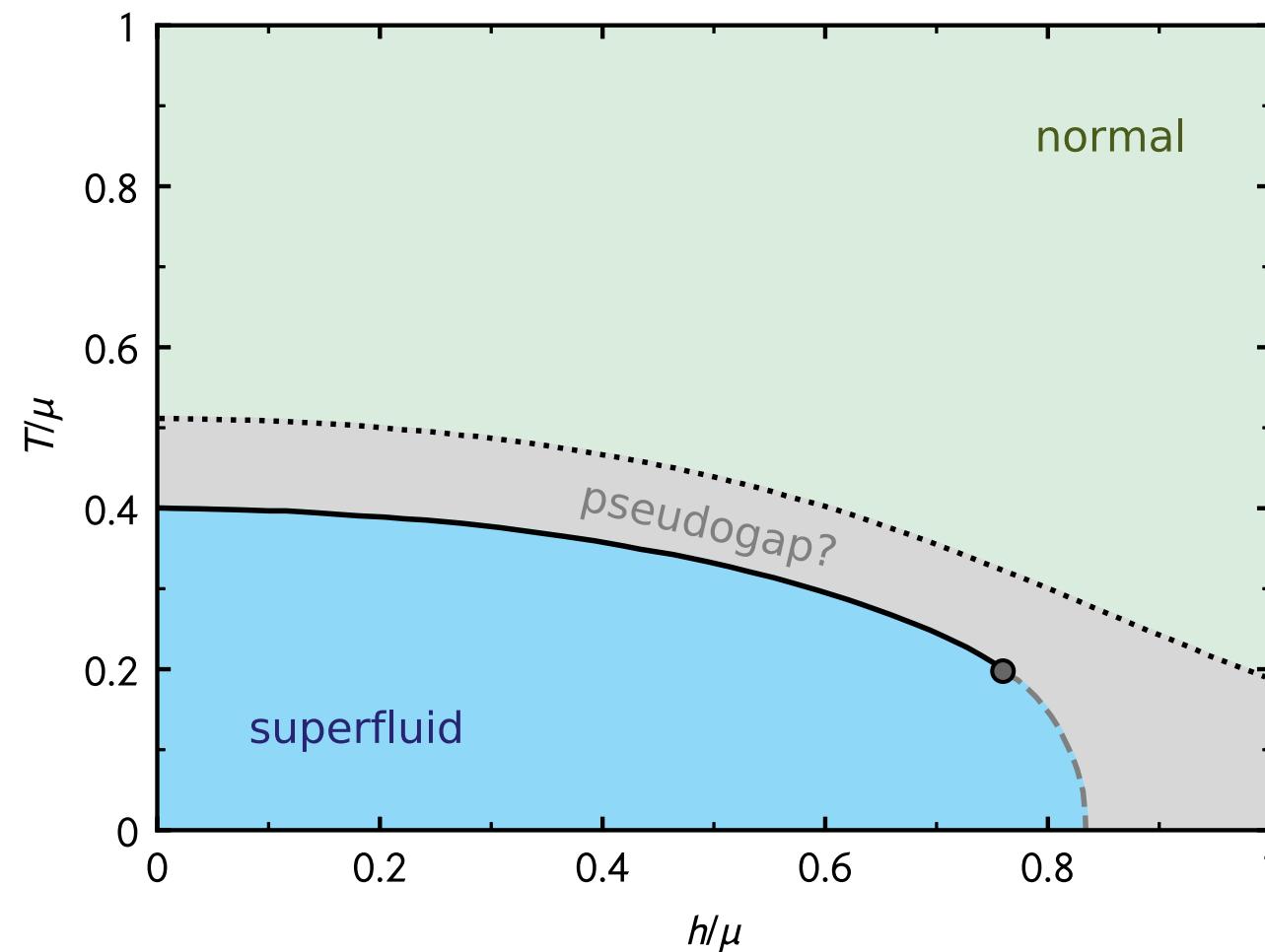
[Scherzer,Seiler,Sexty,Stamatescu '19]

recap: stochastic quantization & CL

- ▶ SQ: interpret **Euclidean field theories** as equilibrium limit of a fictitious **random process**
- ▶ complex Langevin provides a way to **evade sign problems** in some cases
- ▶ however: not guaranteed to work a-priori and the **behavior needs to be monitored carefully**
- ▶ recent review: [Berger et al. arXiv:1907.10183]

the spin-polarized unitary Fermi gas at finite temperature

[LR, Loheac, Drut, Braun '18]



$$\begin{aligned}\mathcal{Z} &= \text{Tr} \left[e^{-\beta(\hat{H}-\mu_{\uparrow}\hat{N}_{\uparrow}-\mu_{\downarrow}\hat{N}_{\downarrow})} \right] \\ &= \text{Tr} \left[e^{-\beta(\hat{H}-\mu\hat{N}-h\hat{M})} \right]\end{aligned}$$

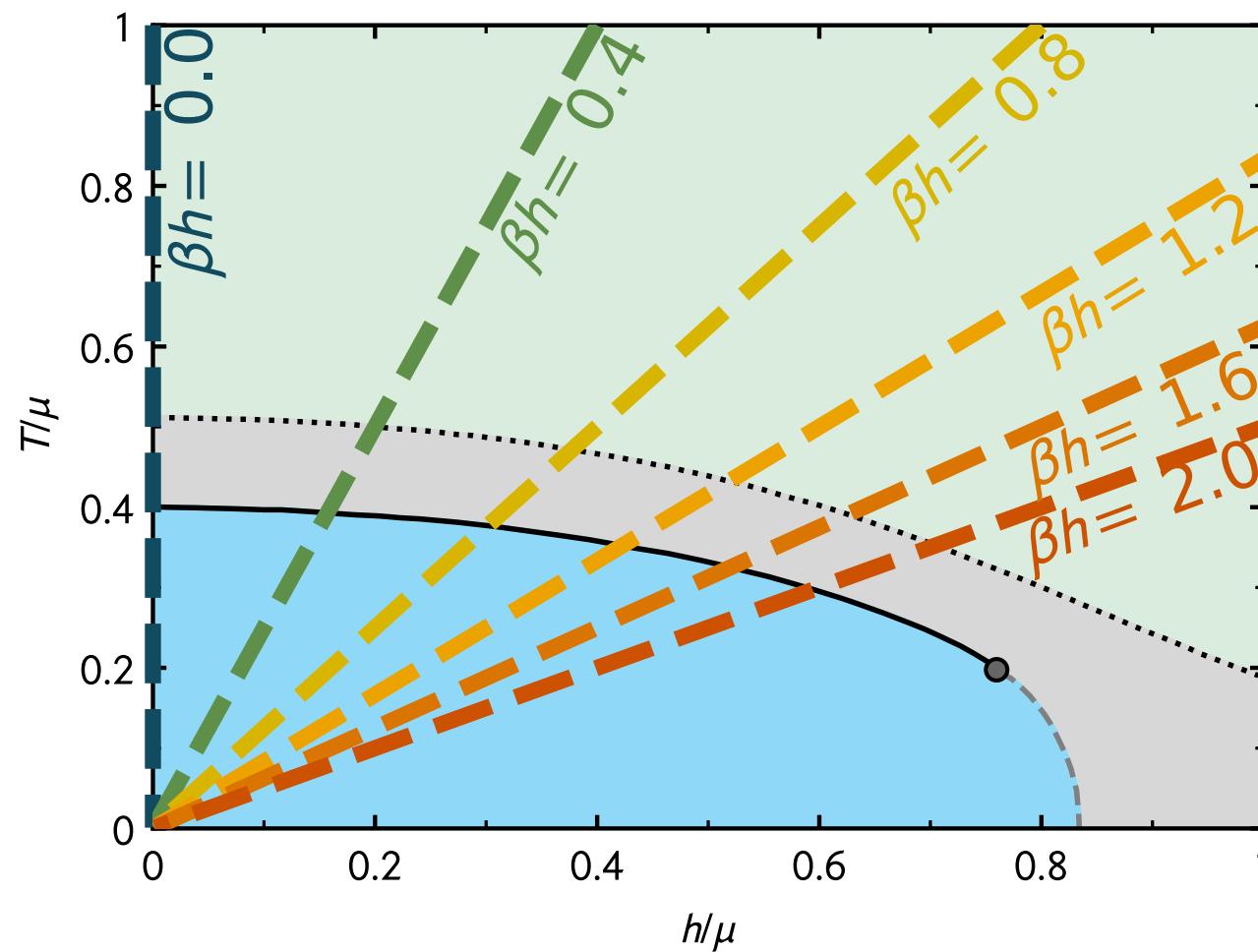
$$\mu = \frac{\mu_{\uparrow} + \mu_{\downarrow}}{2}$$

$$h = \frac{\mu_{\uparrow} - \mu_{\downarrow}}{2}$$

[fRG phase diagram: Boettcher et. al '15]
[LW study of polarized UFG: Frank,Lang,Zwerger '18]

the spin-polarized unitary Fermi gas at finite temperature

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[fRG phase diagram: Boettcher et. al '15]
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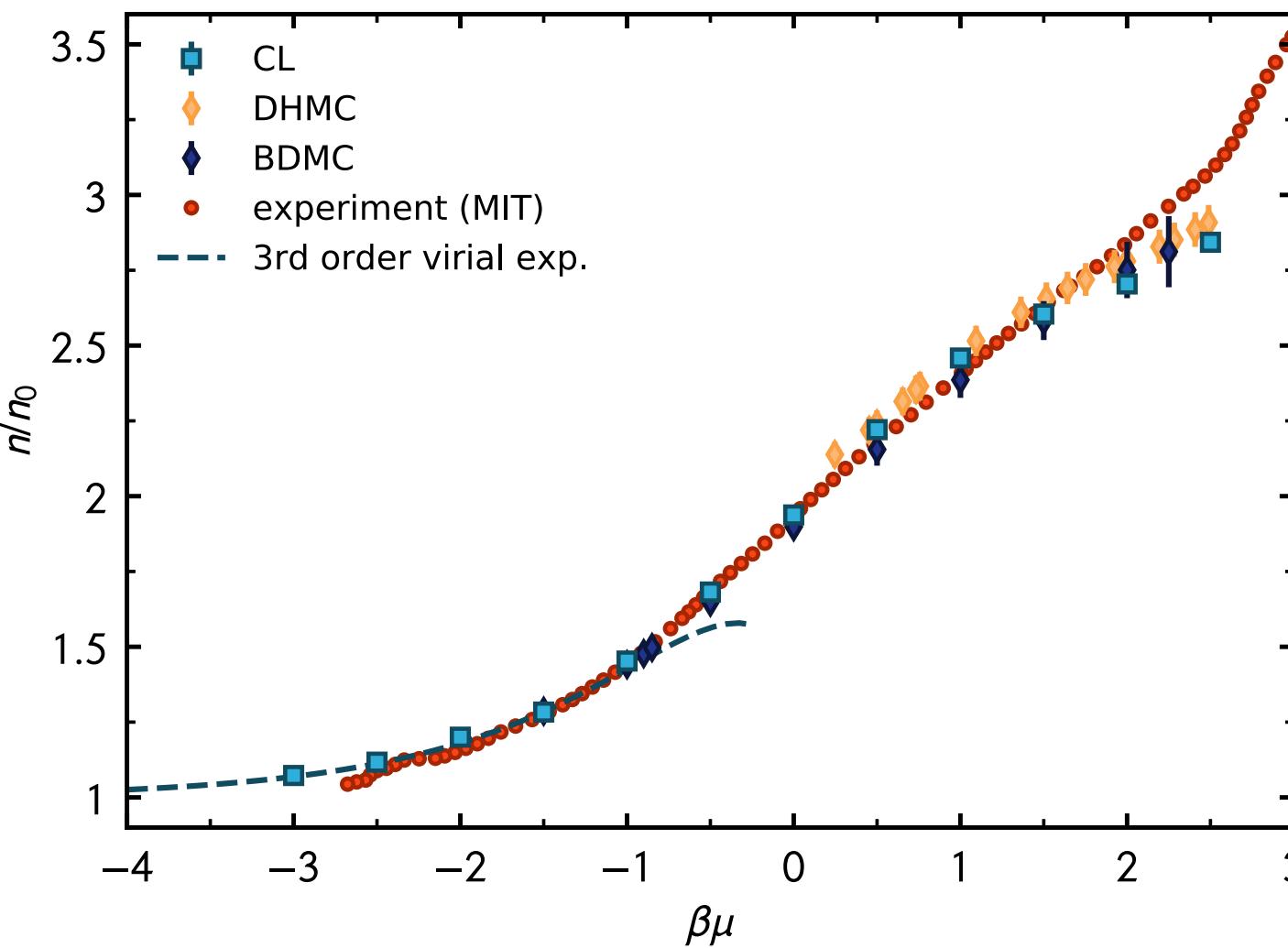
density equation of state

[LR, Loheac, Drut, Braun '18]

[experiment/BDMC: van Houcke et al. '12]

[DHMC: Drut,Lähde,Wlazłowski,Magierski '12]

[diagMC: Rossi,Ohgoe,van Houcke,Werner '18]



classical regime

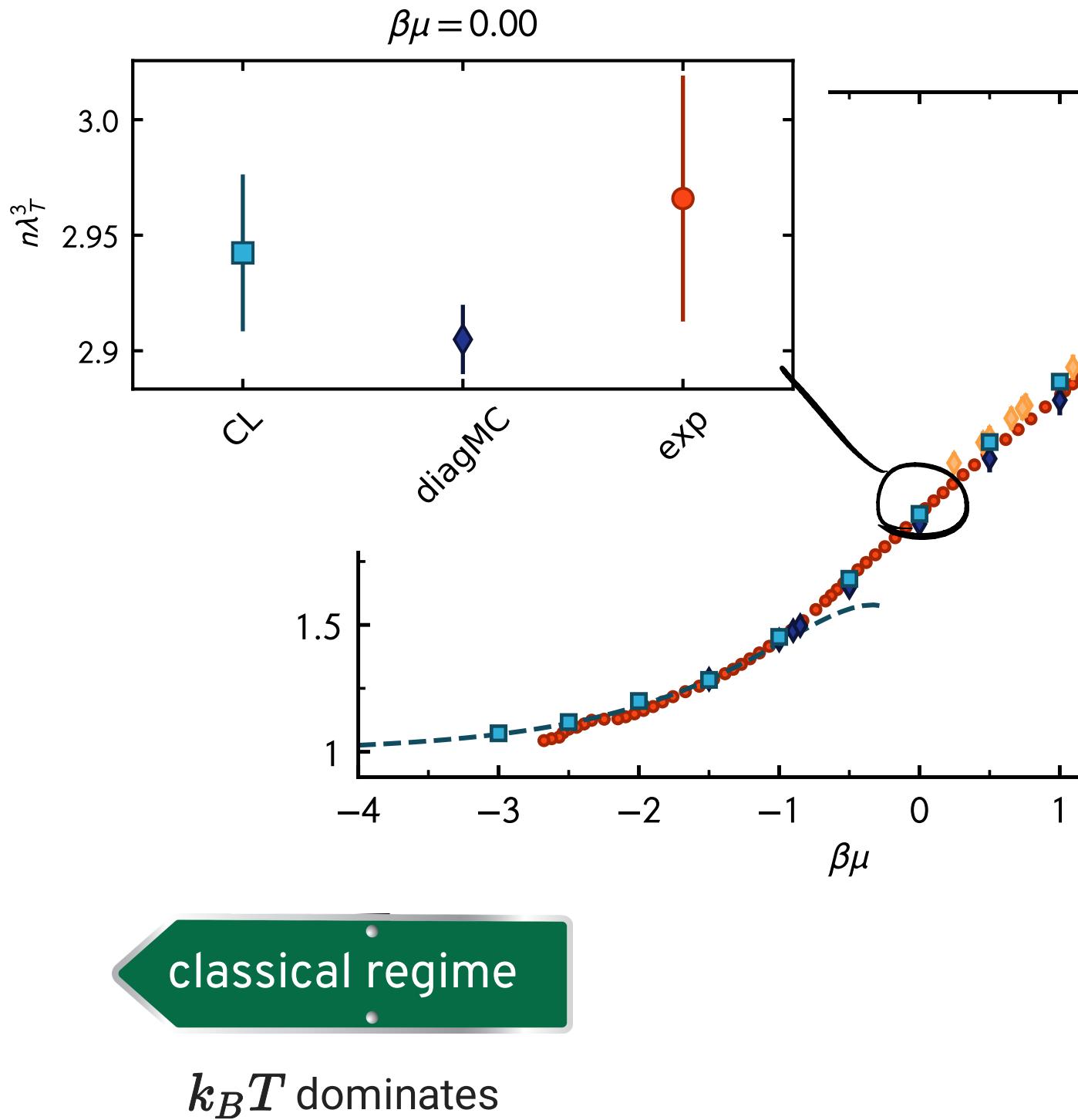
$k_B T$ dominates

quantum regime

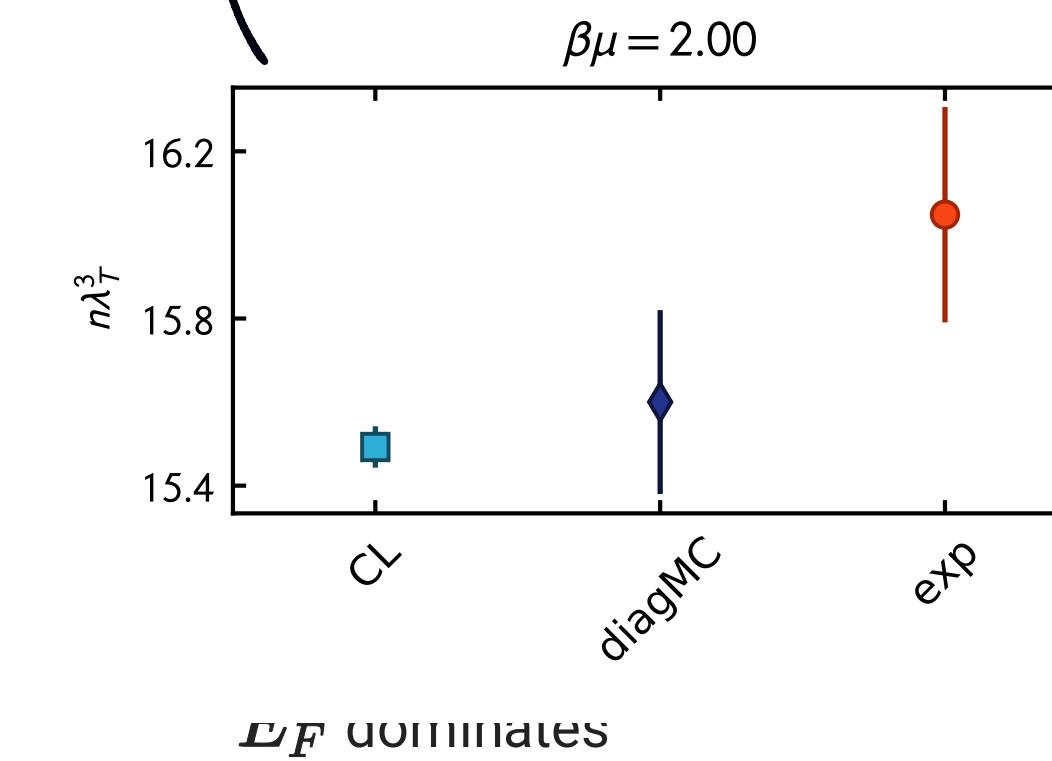
E_F dominates

density equation of state

[LR, Loheac, Drut, Braun '18]



[experiment/BDMC: van Houcke et al. '12]
[DPMC: Drut,Lähde,Wlazłowski,Magierski '12]
[diagMC: Rossi,Ohgoe,van Houcke,Werner '18]



density equation of state

[LR, Loheac, Drut, Braun '18]

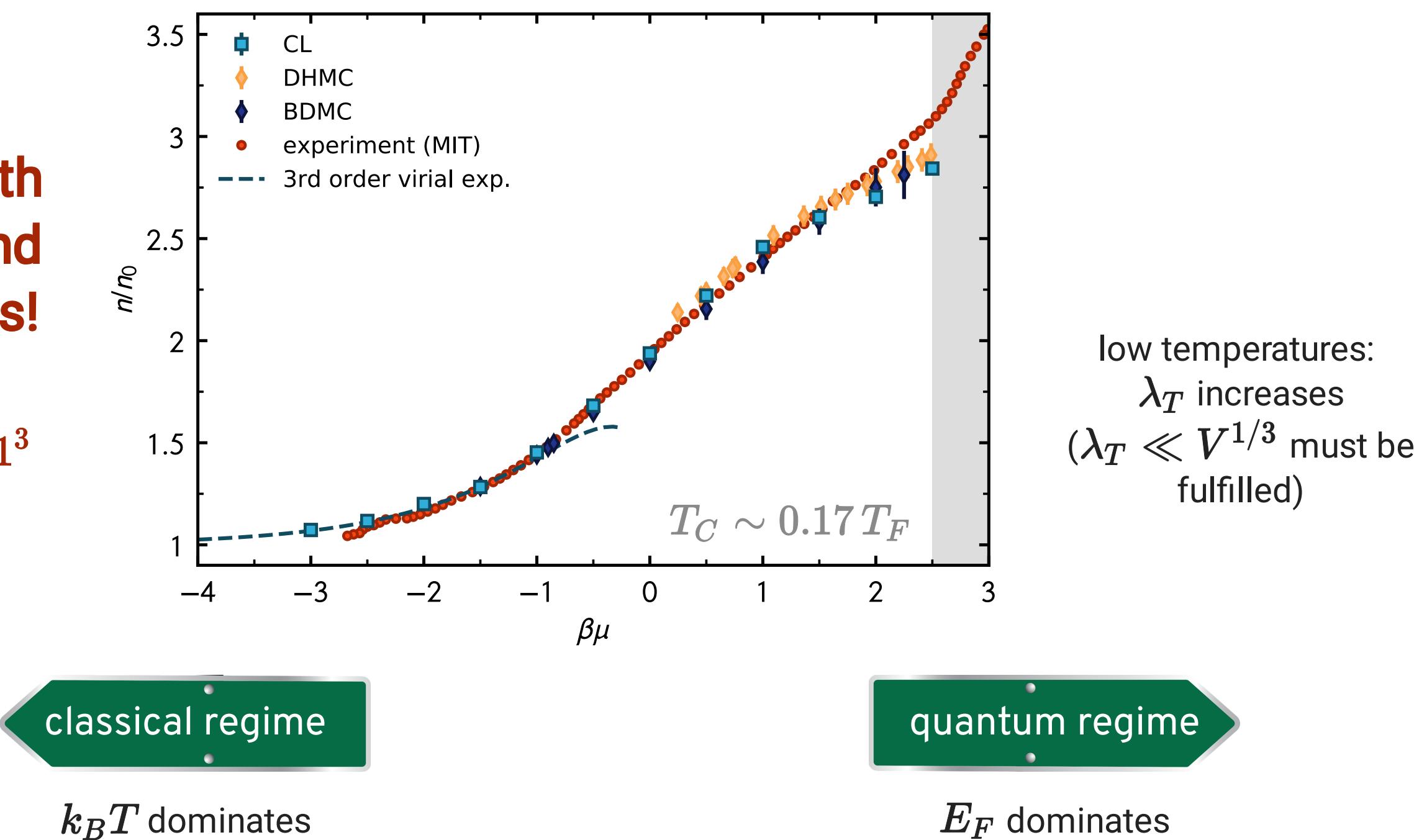
[experiment/BDMC: van Houcke et al. '12]

[DPMC: Drut,Lähde,Wlazłowski,Magierski '12]

[diagMC: Rossi,Ohgoe,van Houcke,Werner '18]

good
agreement with
experiment and
other methods!

CL results:
finite lattice $V = 11^3$



interlude: the virial expansion

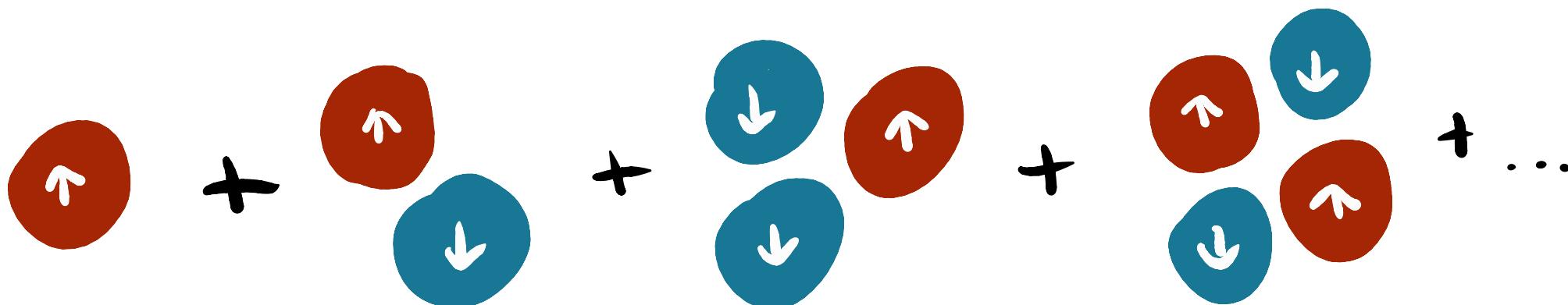
[Liu '13]

dilute gases: few-body correlations dominate

idea: describe the system as **expansion in few-body clusters**

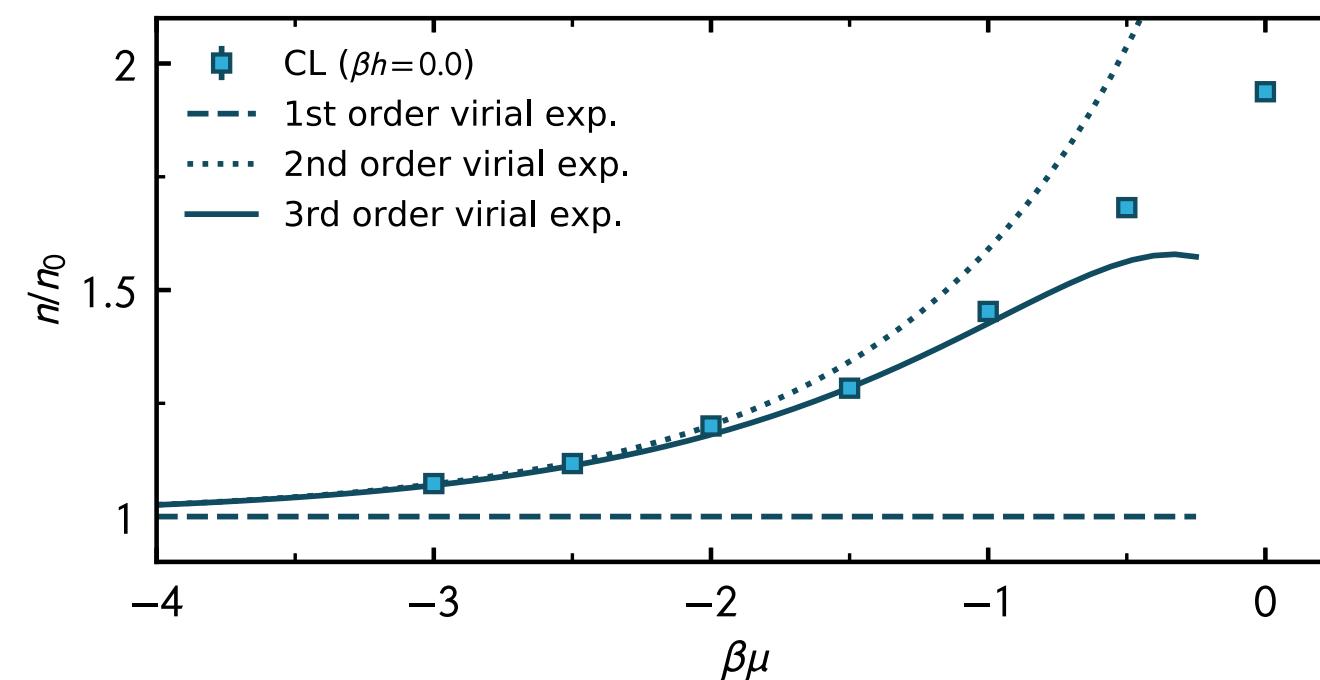
$$z = e^{\beta\mu}$$

$$\ln \mathcal{Z} = Q_1 \sum_n z^n b_n$$



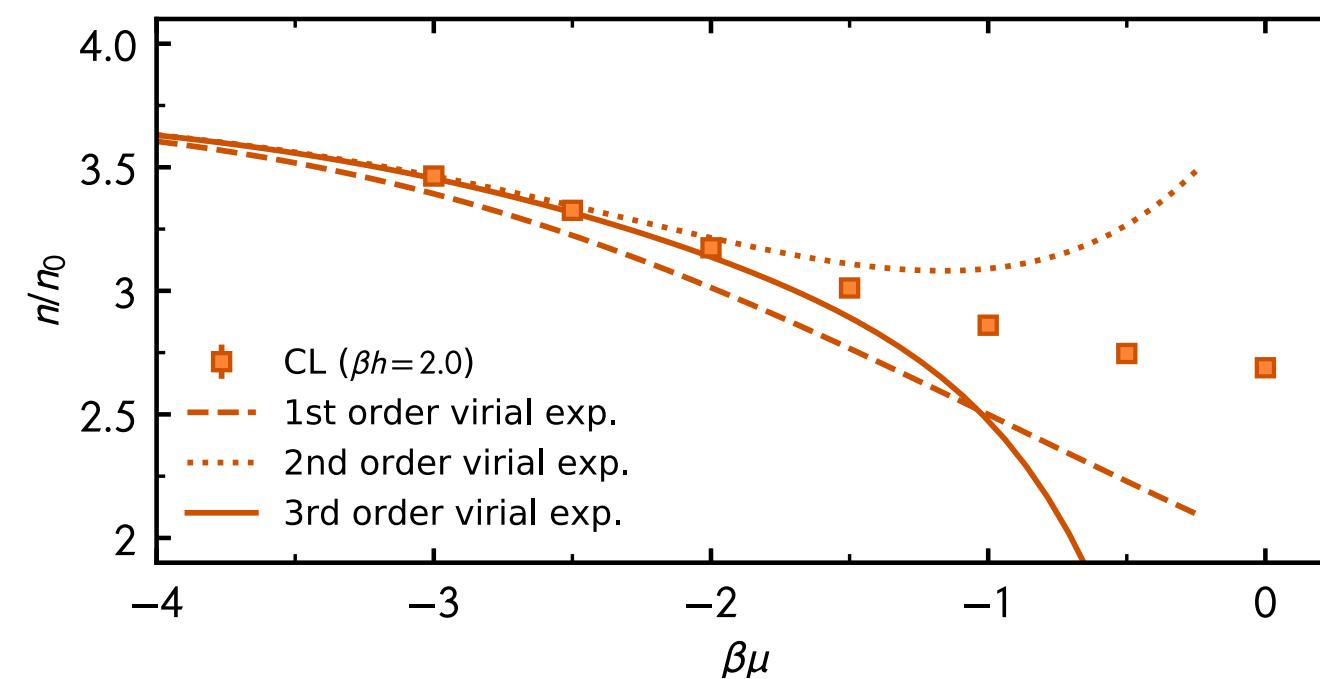
density EOS (virial regime)

[LR, Loheac, Drut, Braun '18]



- ▶ maximum radius of convergence:
 $\beta\mu = 0$

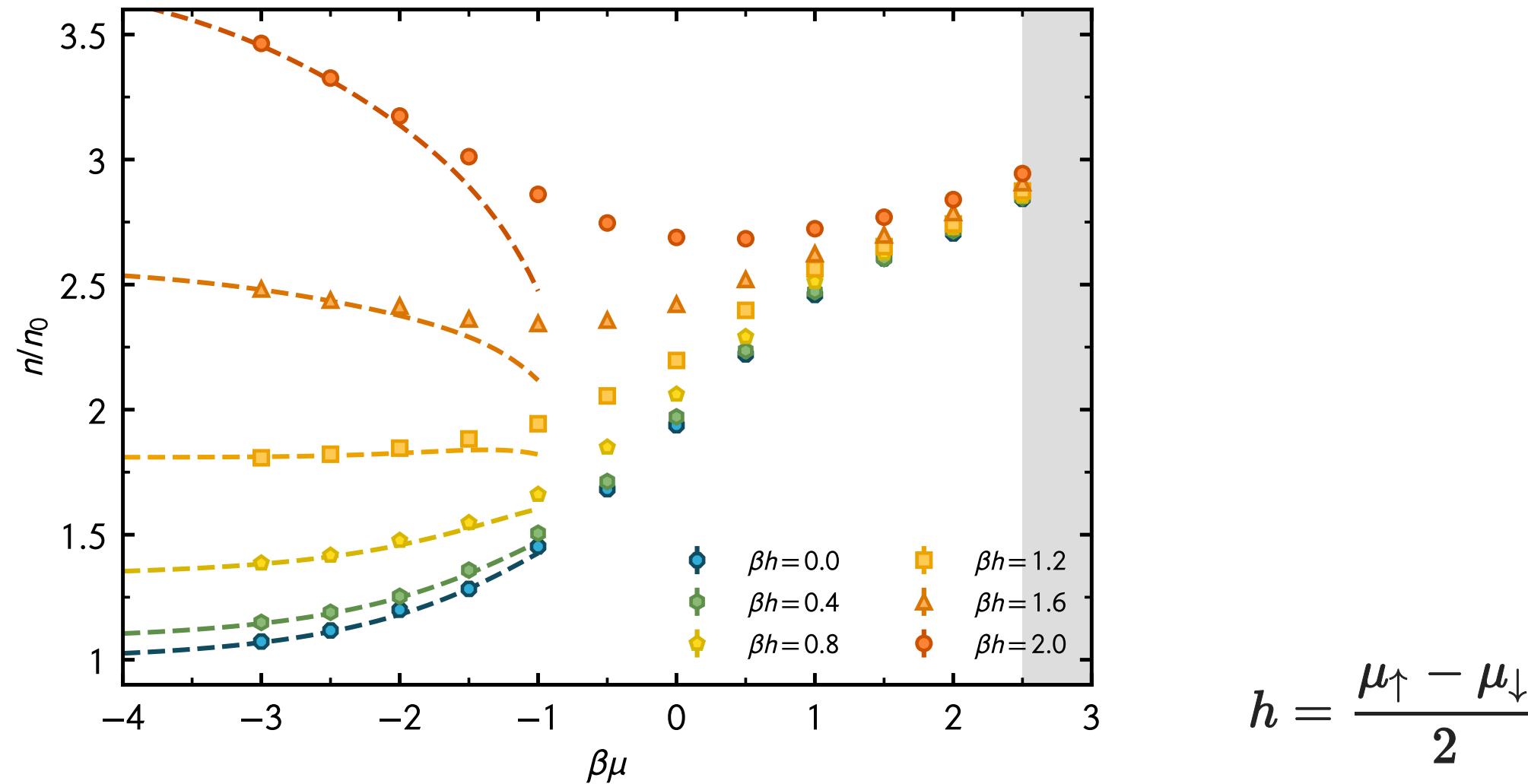
- ▶ **virial expansion approaches the CL results order-by-order**



- ▶ deviation earlier for polarized systems
(majority species hits convergence radius earlier)

density EOS at finite polarization

[LR, Loheac, Drut, Braun '18]



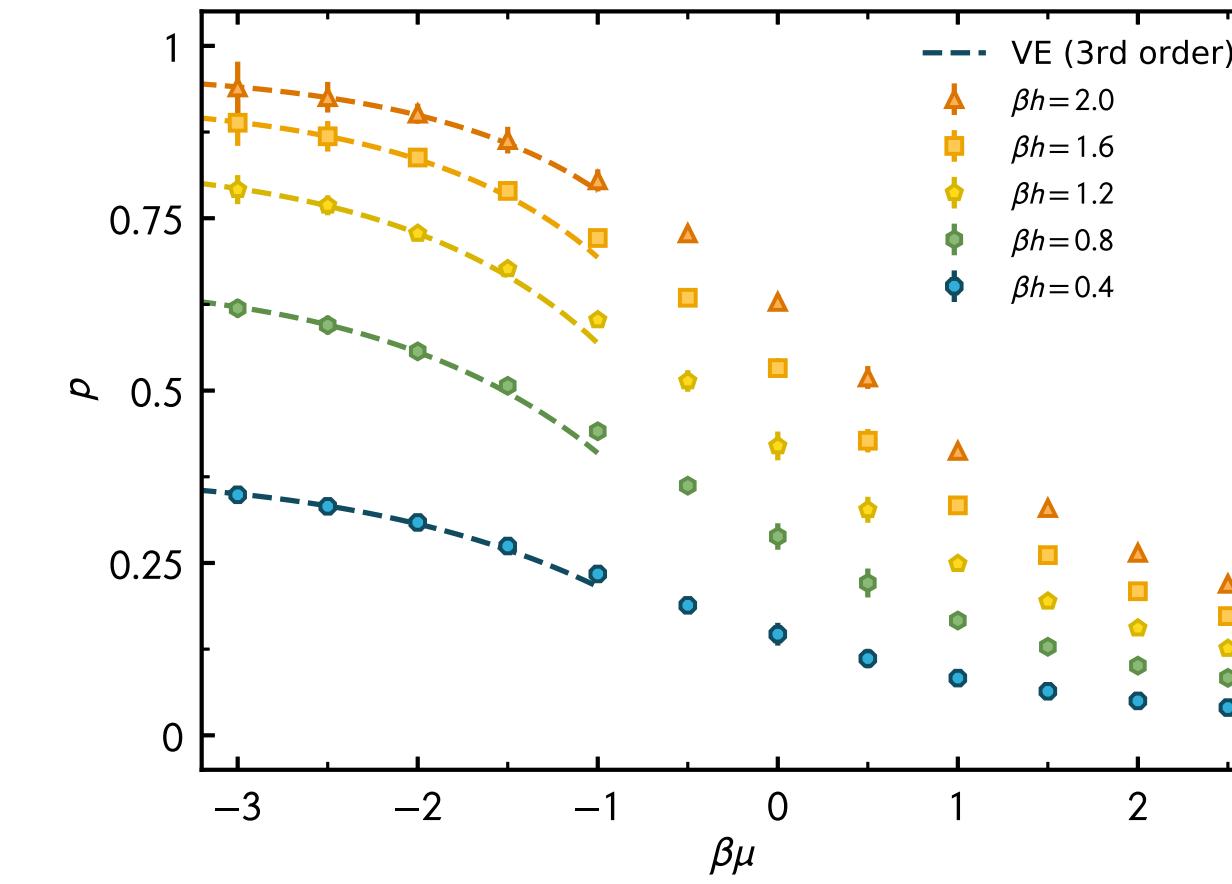
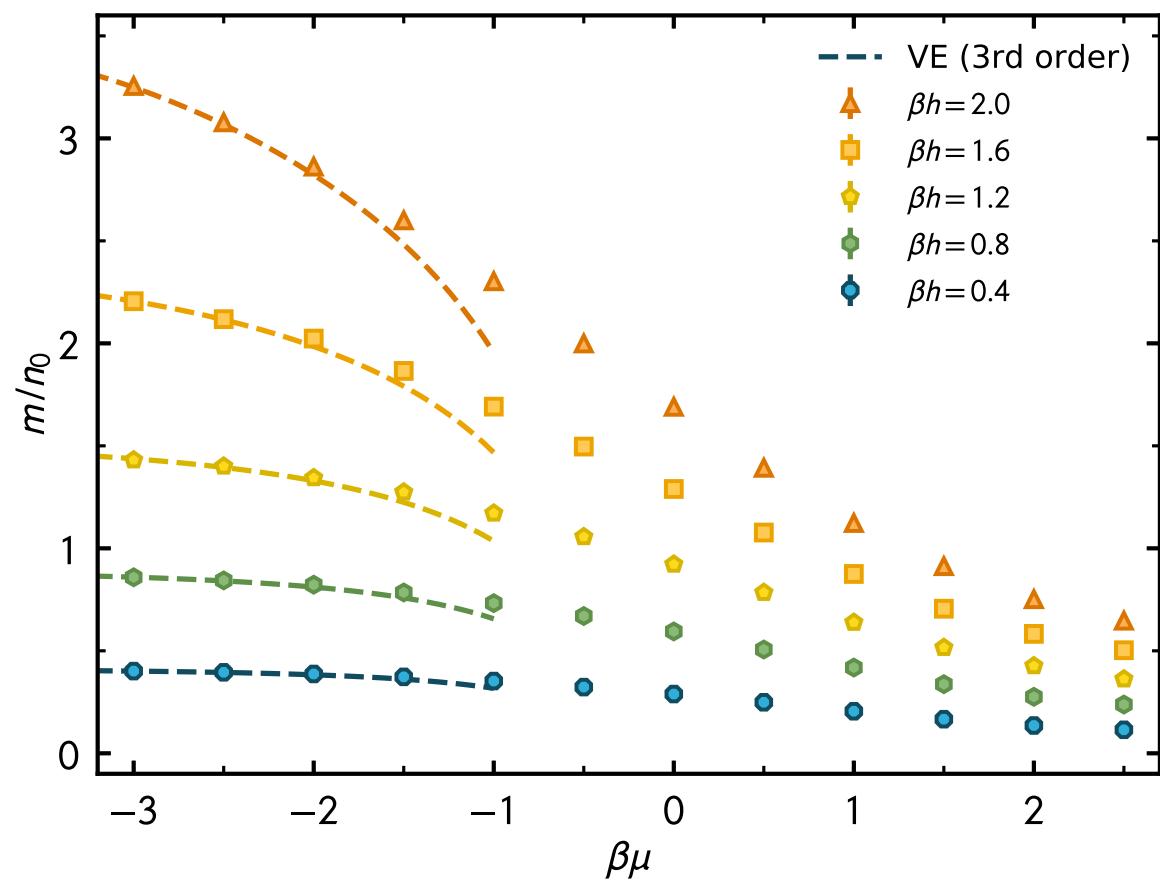
- excellent agreement with virial expansion **for all polarizations**
- **experimentally testable** prediction

magnetization & polarization

[LR, Loheac, Drut, Braun '18]

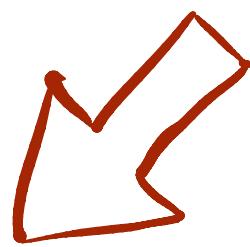
$$m = n_{\uparrow} - n_{\downarrow} = \frac{\partial \ln \mathcal{Z}}{\partial(\beta h)}$$

$$p = \frac{n_{\uparrow} - n_{\downarrow}}{n_{\uparrow} + n_{\downarrow}} = \frac{m}{n}$$



textbook thermodynamics

$$n = \frac{\partial \ln \mathcal{Z}}{\partial(\beta\mu)}$$

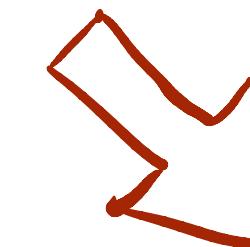


pressure & energy

$$P(\beta\mu) = \frac{1}{\beta} \int_{-\infty}^{\beta\mu} n(x) dx$$

$$E = \frac{3}{2}PV$$

$$m = \frac{\partial \ln \mathcal{Z}}{\partial(\beta h)}$$



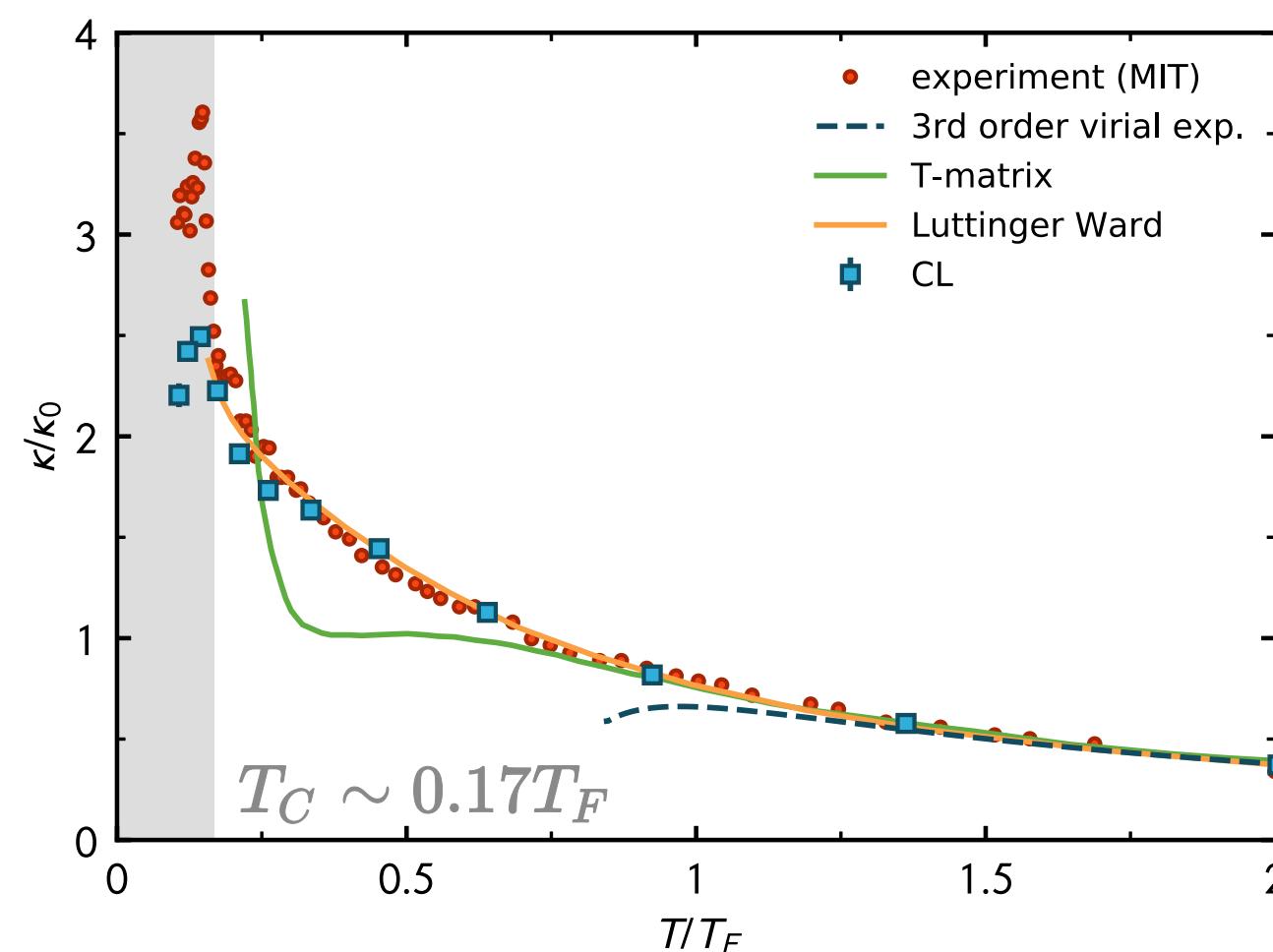
thermodynamic response

$$\kappa = \frac{1}{n} \left(\frac{\partial n}{\partial P} \right)_{T,V,h}$$

$$\chi = \left(\frac{\partial m}{\partial h} \right)_{T,V,\mu}$$

compressibility (unpolarized)

[LR, Loheac, Drut, Braun '18]



$$\kappa = \frac{1}{n} \left(\frac{\partial n}{\partial P} \right)_{T,V,h} = \frac{1}{n^2} \left(\frac{\partial n}{\partial \mu} \right)_{T,V,h}$$

- ▶ sudden increase of κ indicates superfluid phase transition
- ▶ features of curve recovered with CL
- ▶ quantitative disagreement at low temperatures

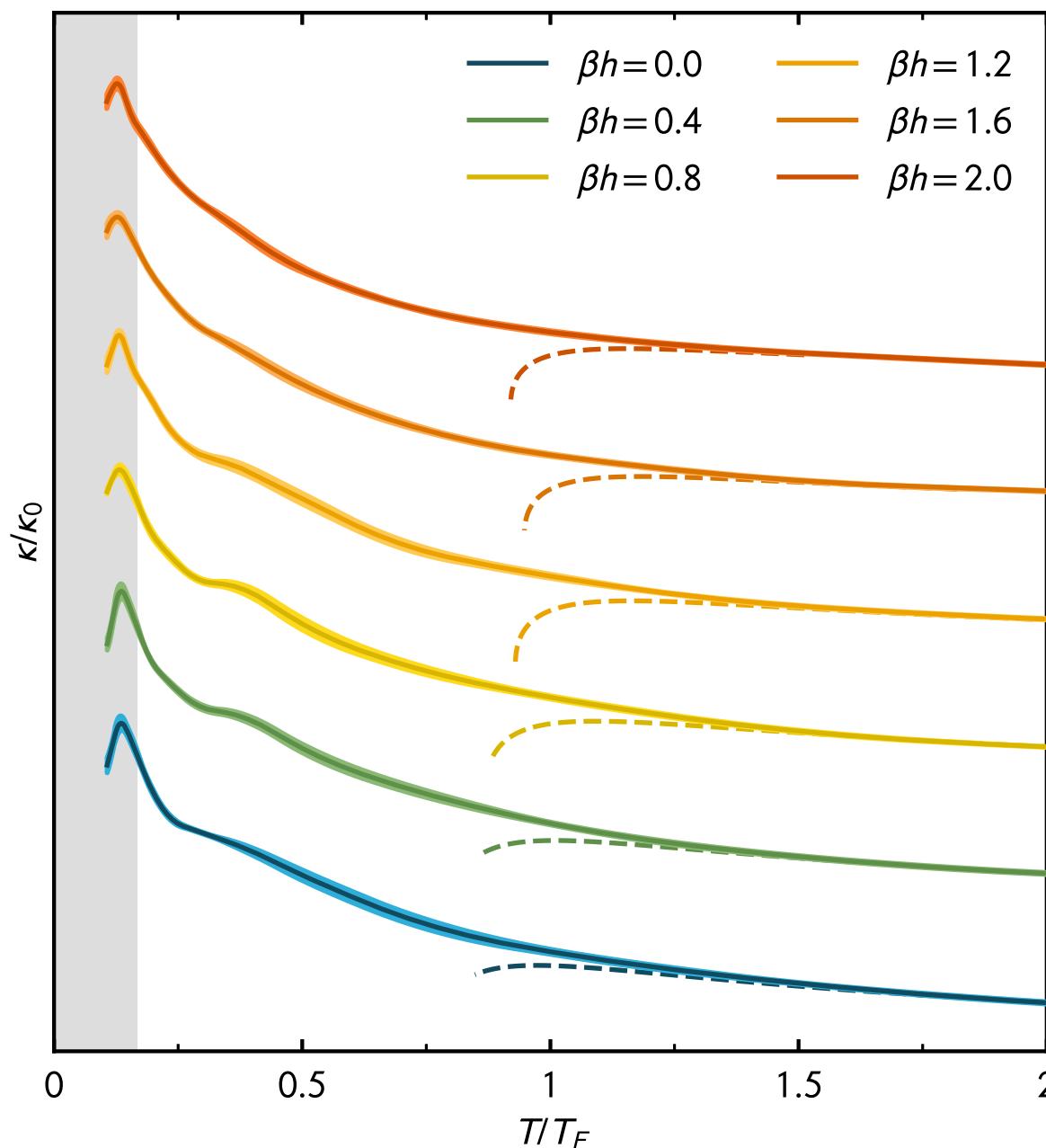
[experiment: Ku,Sommer,Cheuck,Zwierlein '12]

[Luttinger-Ward: Enns,Haussmann '12]

[T-matrix: Pantel et al. '14]

compressibility for polarized systems

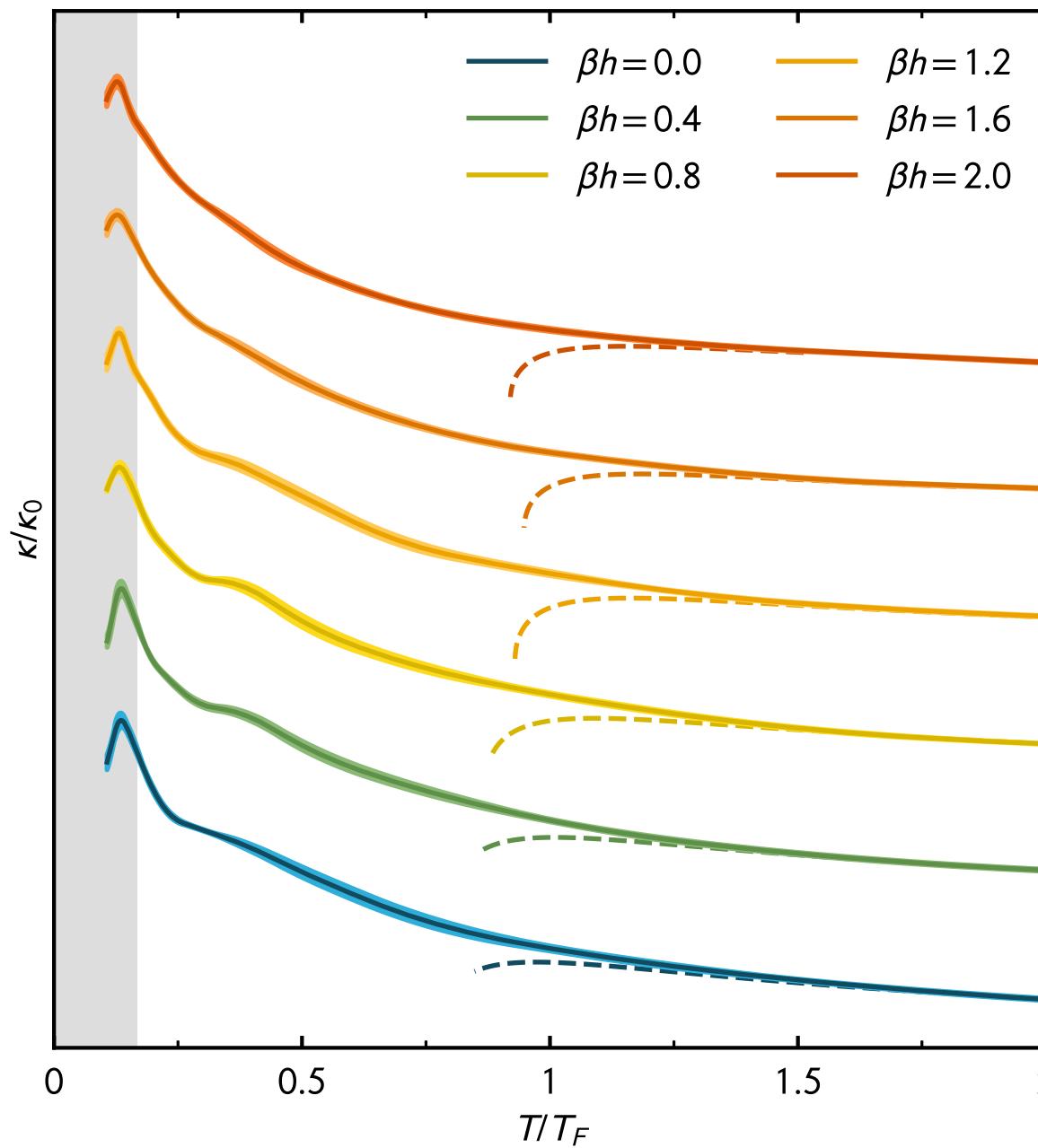
[LR, Loheac, Drut, Braun '18]



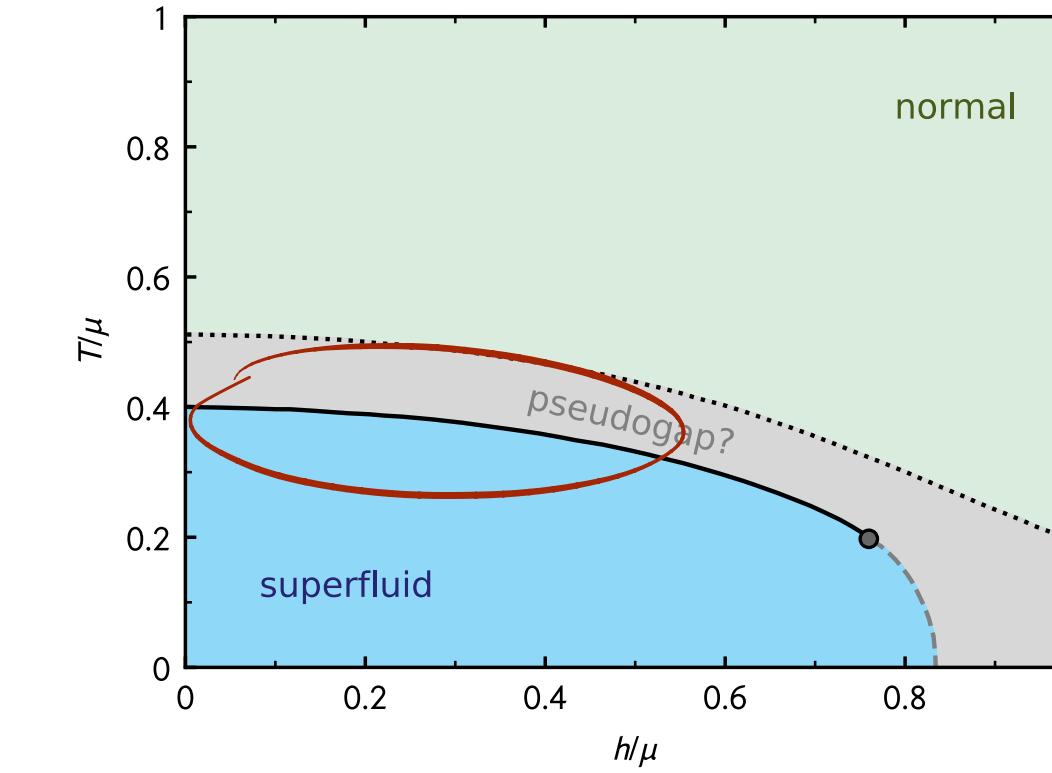
- ▶ weak dependence of the critical temperature on polarization indicated
- ▶ challenging to extract precise T_C

compressibility for polarized systems

[LR, Loheac, Drut, Braun '18]



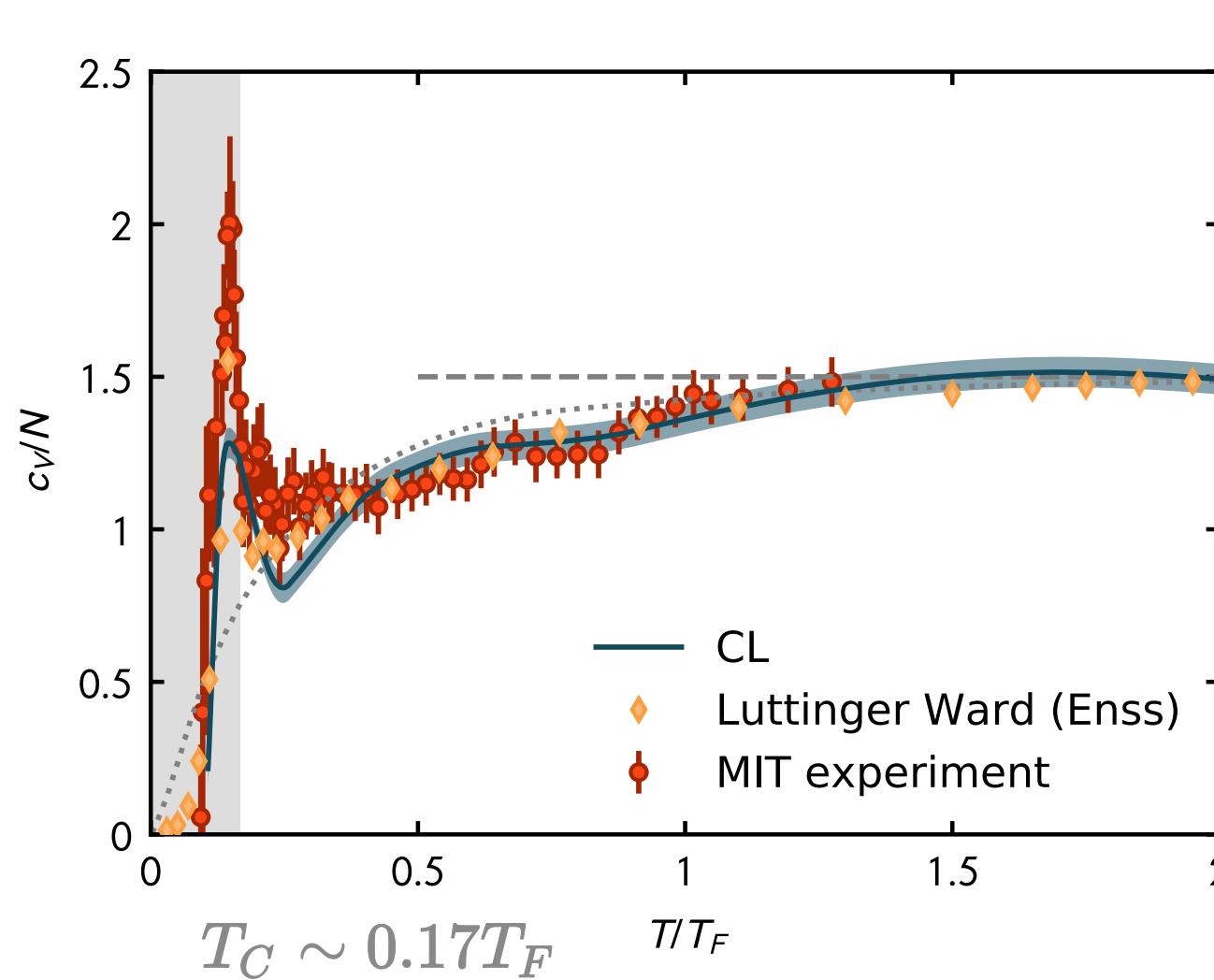
- ▶ weak dependence of the critical temperature on polarization indicated
- ▶ challenging to extract precise T_C



[fRG PD: Boettcher et. al '15]

specific heat (unpolarized)

[LR, Drut, Braun in preparation]



[experiment: Ku,Sommer,Cheuck,Zwierlein '12]
[Luttinger-Ward: Enns,Haussmann '12]

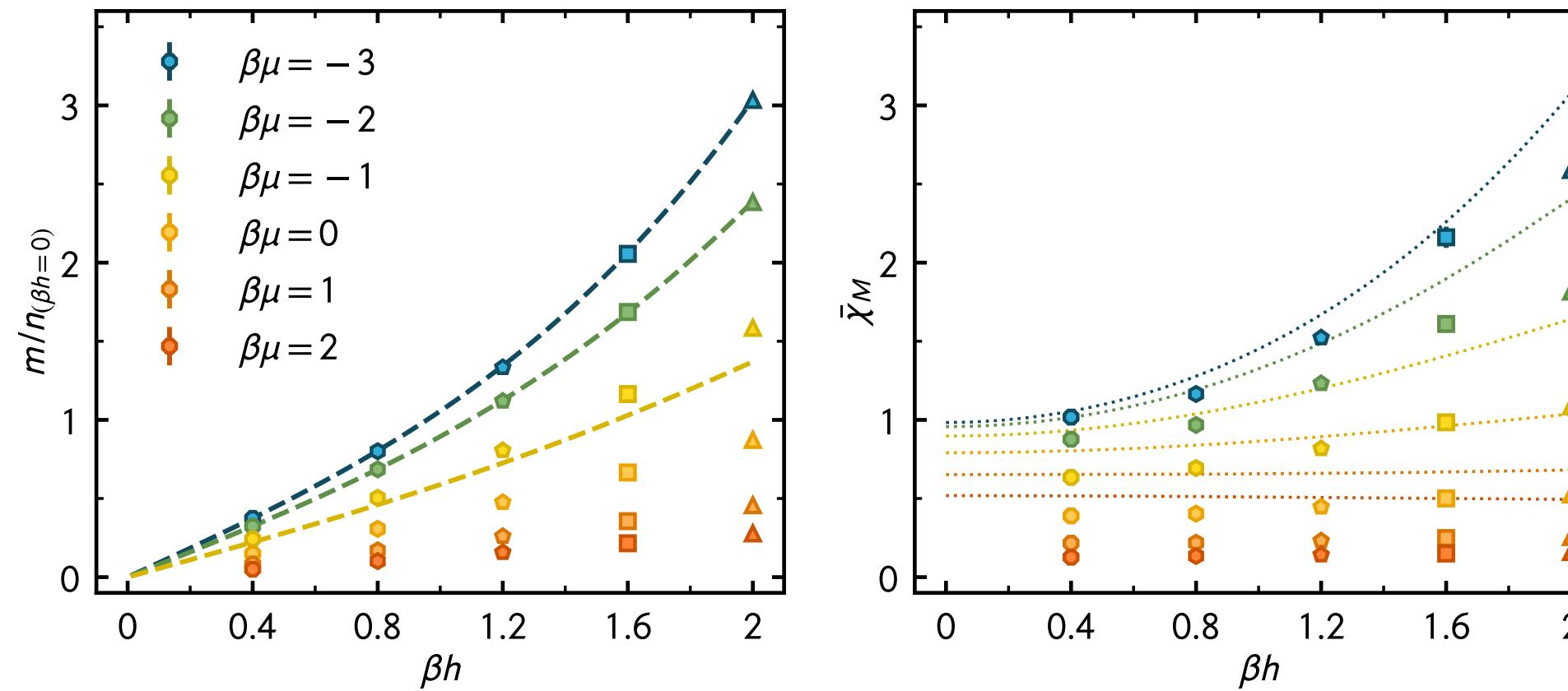
$$\frac{C_V}{N} = \frac{1}{N} \left(\frac{\partial E}{\partial T} \right)_{N,V}$$

- ▶ high temperatures: Dulong-Petit law
 $C_V/N = \frac{3}{2}$
- ▶ second order transition at T_C ("lambda transition")
- ▶ **preliminary data (polarized values not yet available)**

spin susceptibility

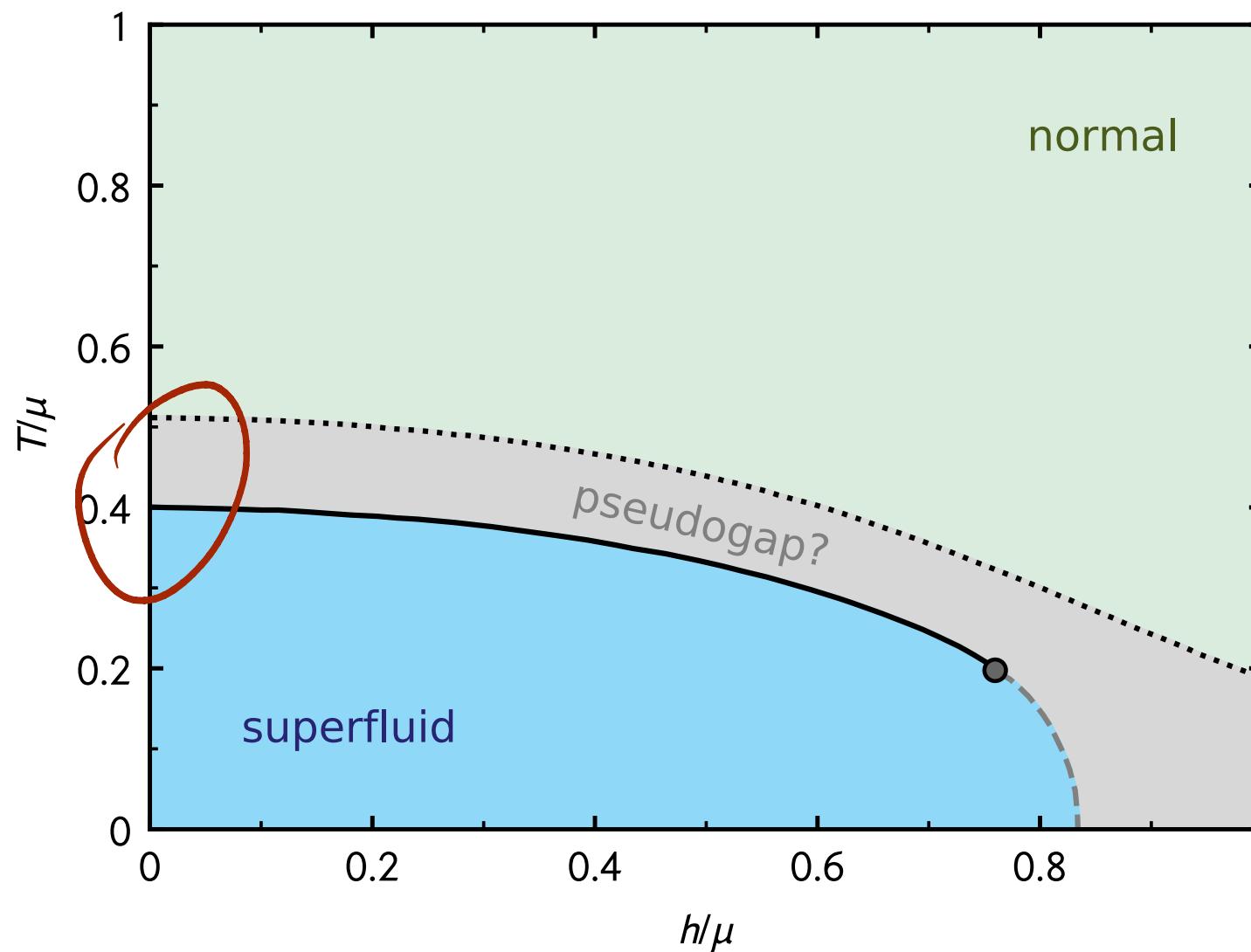
[LR, Loheac, Drut, Braun '18]

$$\chi = \left(\frac{\partial m}{\partial h} \right)_{T,V,\mu}$$



- ▶ Pauli susceptibility field independent at low field and temperature
- ▶ **UFG: dependence on βh very similar to FG, but rescaled**

UFG phase diagram (sketch)



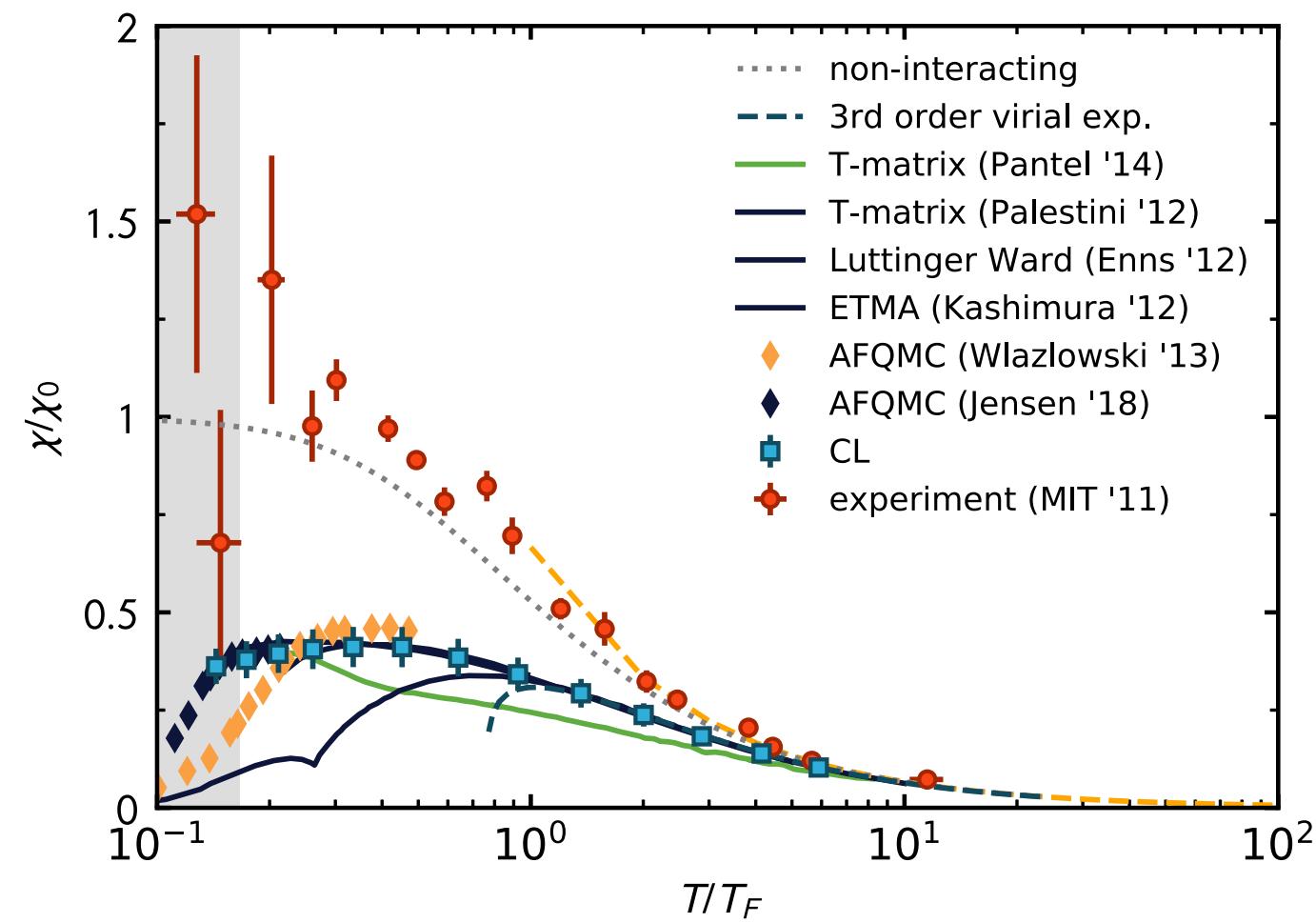
$$\mu = \frac{\mu_\uparrow + \mu_\downarrow}{2}$$

$$h = \frac{\mu_\uparrow - \mu_\downarrow}{2}$$

[fRG PD: Boettcher et. al '15]

magnetic susceptibility (unpolarized)

[LR, Loheac, Drut, Braun in preparation]



[recent review: Jensen et al. '18]

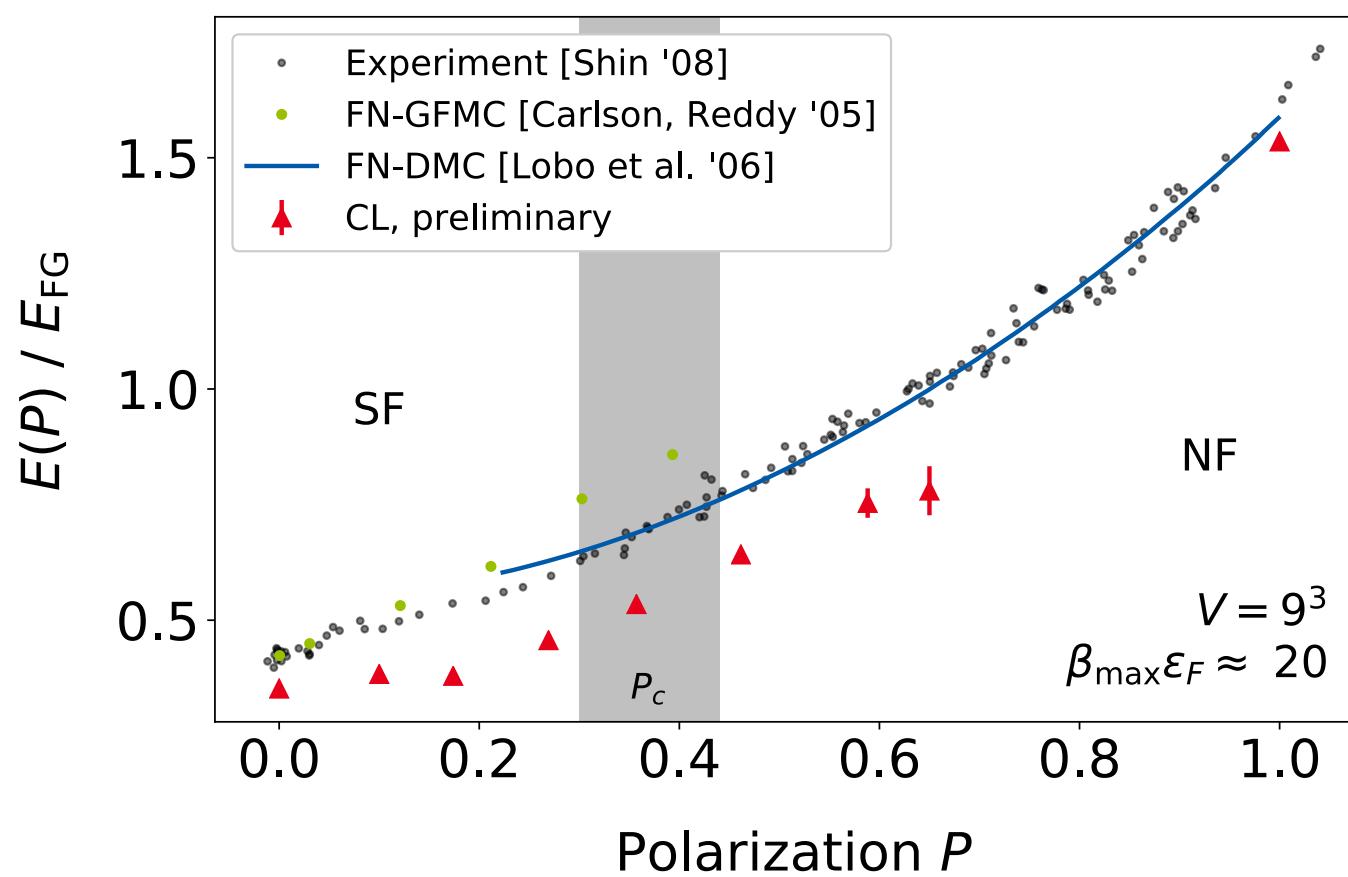
- ▶ high temperature:
Curie's law $\chi \propto T^{-1}$
- ▶ theory & experiment agree at high temperatures
- ▶ Pseudogap: suppression of χ at $T > T_C$
- ▶ **low temperature: discrepancy between experiment and theory**
- ▶ **CL: pseudogap possible T^* and T_C seem to be very close**

recap: unitary fermions

- ▶ **EOS, magnetic properties & response**
accessible for the unitary Fermi gas at finite temperature and polarization
- ▶ CL **matches state-of-the art results** from other methods and experiments wherever available

preliminary: UFG in the ground state

[Ehmann, LR, Drut, Braun in preparation]

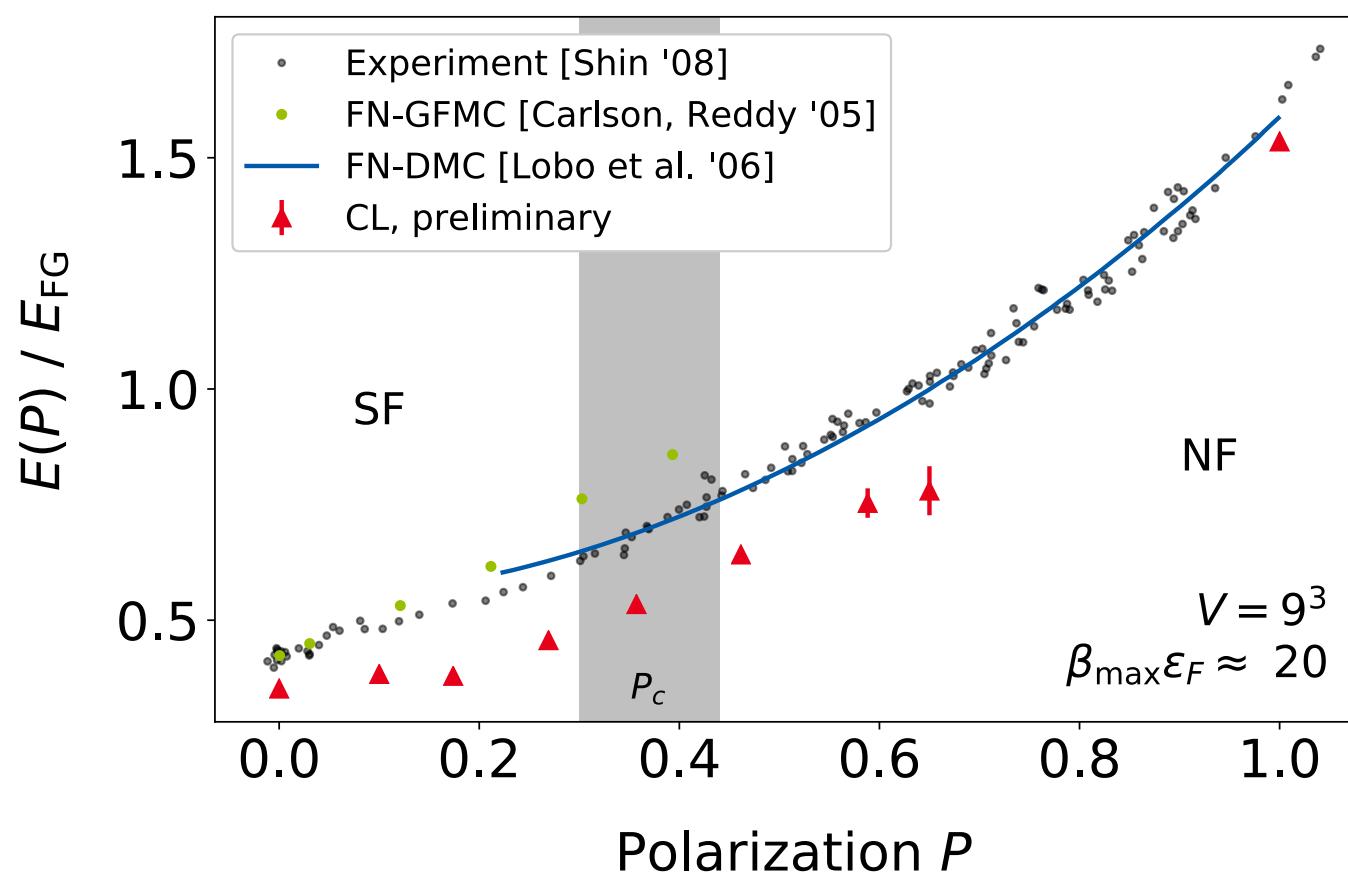


- ▶ qualitative agreement of preliminary CL data with other MC values & experiment
- ▶ **CL does not require an ansatz**
- ▶ no continuum extrapolations as of yet
- ▶ Outlook: **FFLO ground state** near critical polarization?

[experiment: Shin '08]
[FN_GFMC: Carlson,Reddy '05]
[FN-DMC: Lobo et al. '06]

preliminary: UFG in the ground state

[Ehmann, LR, Drut, Braun in preparation]

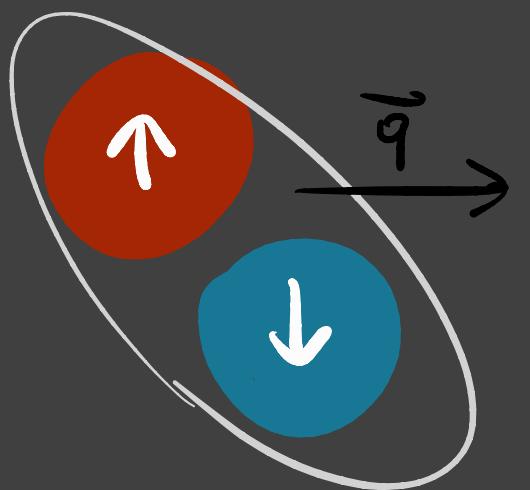


[experiment: Shin '08]
[FN_GFMC: Carlson,Reddy '05]
[FN-DMC: Lobo et al. '06]

- ▶ qualitative agreement of preliminary CL data with other MC values & experiment
- ▶ **CL does not require an ansatz**
- ▶ no continuum extrapolations as of yet
- ▶ Outlook: **FFLO ground state** near critical polarization?

Check out the poster of Florian Ehmann for more details!

STAY TUNED!

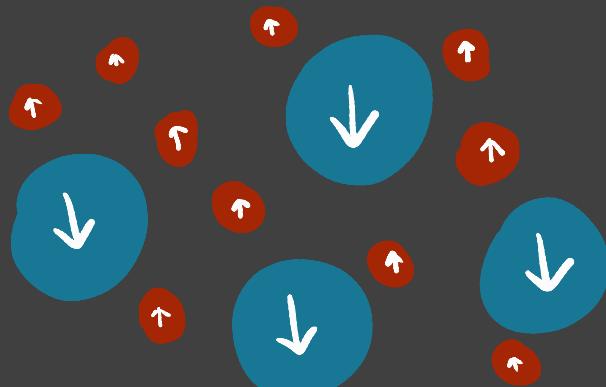


looking for inhomogeneous phases in the UFG

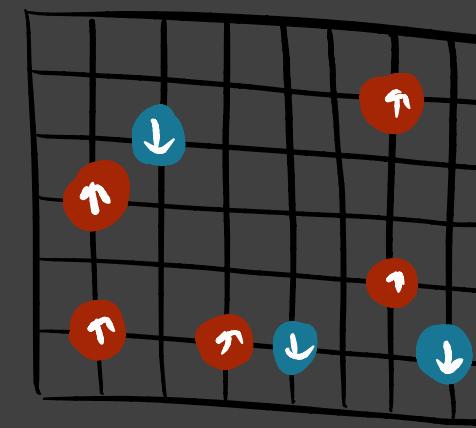
[ongoing work with Florian Ehmann, Joaquin Drut & Jens Braun]

thermodynamics of 2D fermions at finite polarization

[ongoing work with Josh McKenney, Andrew Loheac, Joaquin Drut & Jens Braun]



effect of mass-imbalance on fermion pair formation

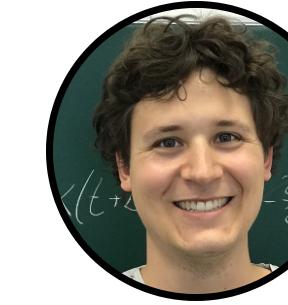


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