



Towards in-medium heavy quarkonium dynamics from first principles

Alexander Rothkopf

Faculty of Science and Technology Department of Mathematics and Physics University of Stavanger

References:

T. Miura, Y. Akamatsu, M. Asakawa, A.R. arXiv:1908.06293
P. Petreczky, A.R., J. Weber, NPA982 (2019) 735
S. Kim, P. Petreczky, A.R. JHEP 1811 (2018) 088
B. Krouppa, A.R., M. Strickland PRD97 (2018) 016017
D. Lafferty, A.R., arXiv:1906.00035, A. Lehmann, A.R. (in preparation)



THE 1ST INTERNATIONAL CONFERENCE ON QUANTUM SYSTEMS IN EXTREME CONDITIONS- SEPTEMBER 24TH 2019 – HD

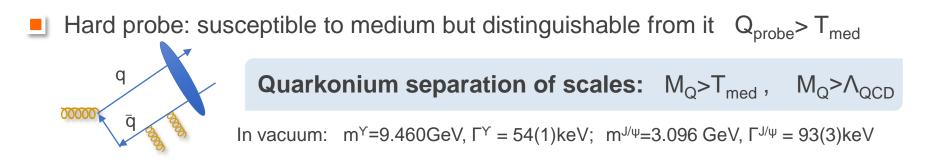
Impurity physics of the early universe

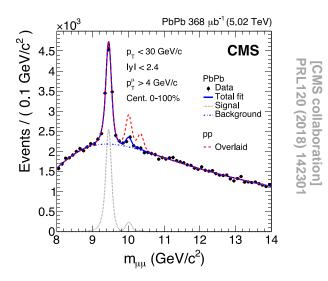


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Bottomonium: a non-equilibrium probe of the full QGP evolution

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OWARDS IN-MEDIUM HEAVY QUARKONIUM DYNAMICS FROM FIRST PRINCIPLES

PbPb 368 µb⁻¹ (5.02 TeV)

 Data Total fit

Signal

- Background

Overlaid

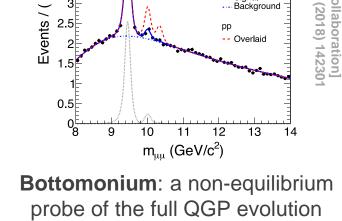
CMS

PRL120

CMS collaboration]

Impurity physics of the early universe

Hard probe: susceptible to medium but distinguishable from it $Q_{probe} > I_{med}$ q **Quarkonium separation of scales:** $M_Q > T_{med}$, $M_Q > \Lambda_{QCD}$ 00000 In vacuum: $m^{\gamma}=9.460$ GeV, $\Gamma^{\gamma}=54(1)$ keV; $m^{J/\psi}=3.096$ GeV, $\Gamma^{J/\psi}=93(3)$ keV



< 30 GeV/c

|y| < 2.4

p^µ₊ > 4 GeV/c

Cent. 0-100%

×10³

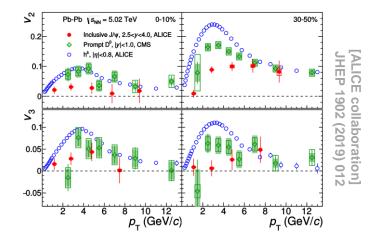
4.5

3.5

3

2.5

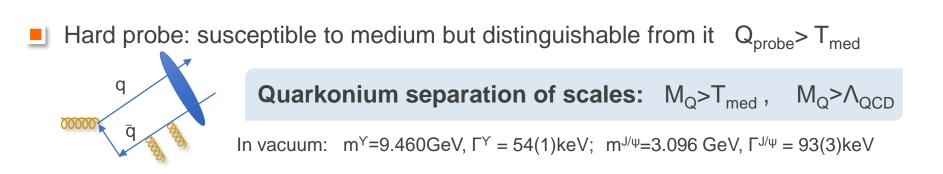
Events / (0.1 GeV/c²)

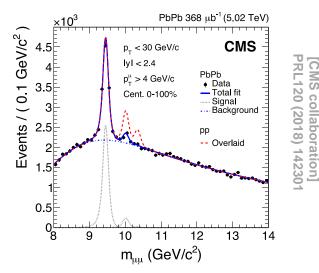


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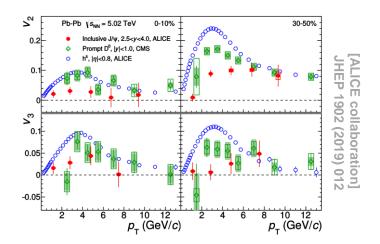
Charmonium: a partially equilibrated probe, sensitive to the late stages

Impurity physics of the early universe





Bottomonium: a non-equilibrium probe of the full QGP evolution



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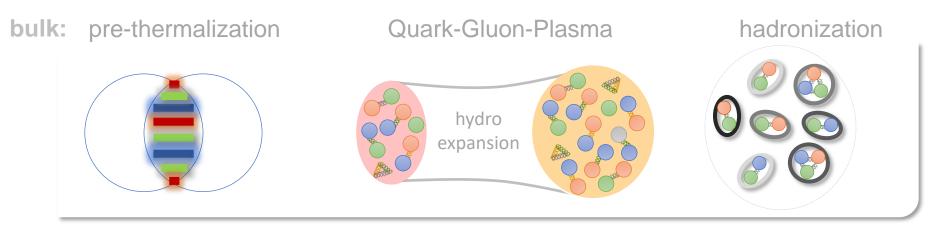
Charmonium: a partially equilibrated probe, sensitive to the late stages

Goal: provide **first principles interpretation** to intricate phenomenology

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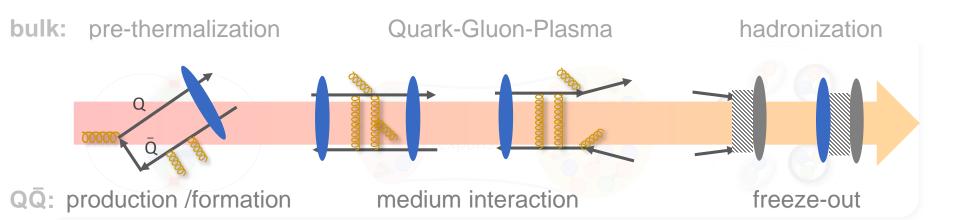
Open theory questions





Open theory questions

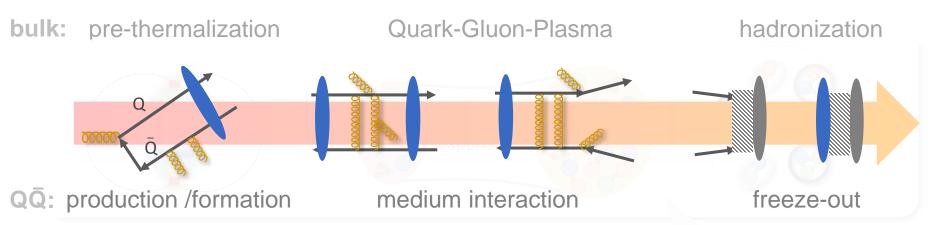
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Open theory questions

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Properties of equilibrium $Q\bar{Q}$

First principles extraction of the heavy quark potential

Novel phenomenological definition of the **Debye mass**

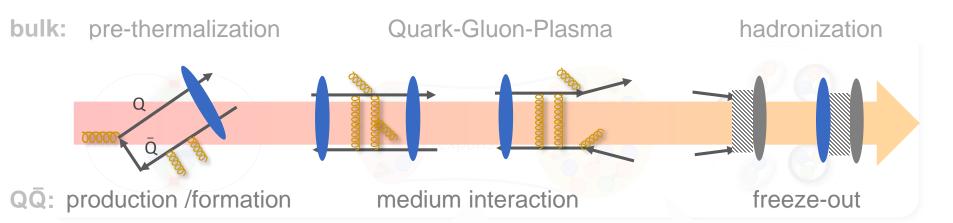
Extraction of thermal spectral properties on the lattice

with P. Petreczky, J. Weber: NPA982 (2019) 735 S. Kim, P. Petreczky, A.R. JHEP 1811 (2018) 088 with D. Lafferty arXiv:1906.00035

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Real-time QQ evol. in local thermal equilibrium

Beyond Schrödinger:

Open-quantum-systems descr. of real-time evolution

Connecting OQS to EFT language of potential

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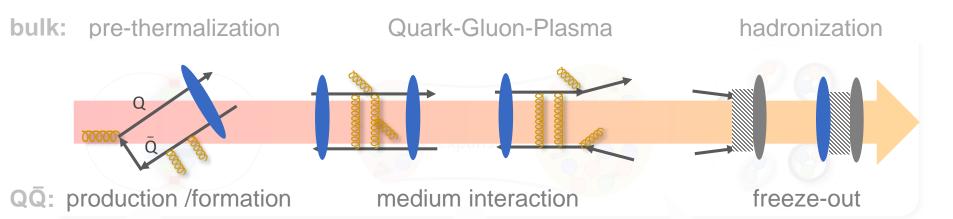
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$Q\bar{Q}$ realtime evol. in the initial stages

First exploratory steps in the glasma

classical statistical simulations for gluons & real-time NRQCD

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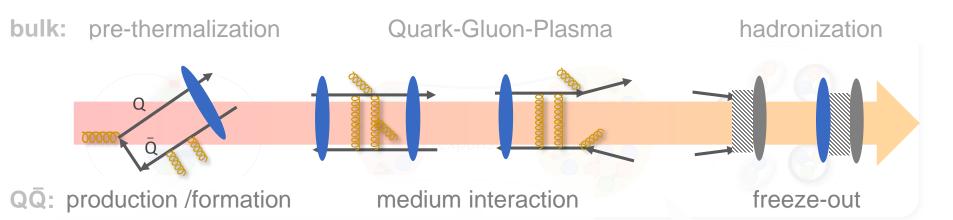
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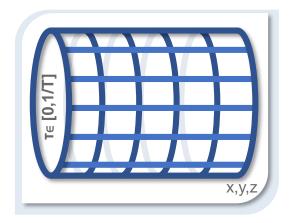
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A robust tool: lattice QCD



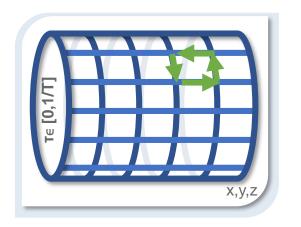
Non-perturbative 1st principles approach to Quantum Chromo Dynamics



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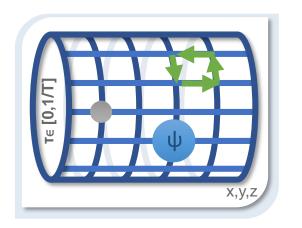


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A robust tool: lattice QCD

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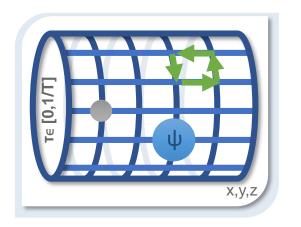


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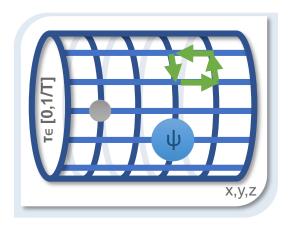
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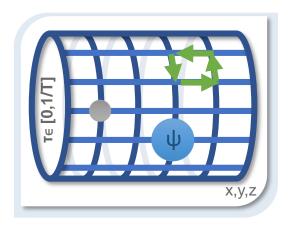
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- **I** Finite extend in **imaginary** time: $1/T = \beta = N_T a_T$

$$\langle O(\mathbf{U}) \rangle = \int \mathcal{D}\mathbf{U} O(\mathbf{U}) e^{-S_{\rm E}^{\rm QCD}[\mathbf{U}]}$$

A robust tool: lattice QCD

Non-perturbative 1st principles approach to Quantum Chromo Dynamics



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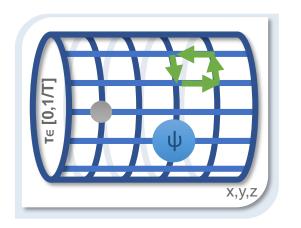
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Finite extend in **imaginary** time:
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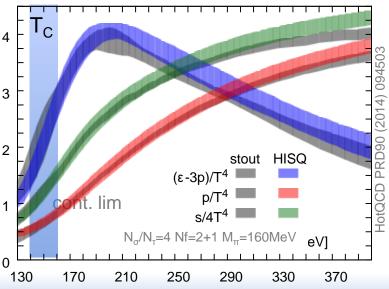
 $\langle O \rangle = \frac{1}{N} \lim_{N \to \infty} \sum_{k=1}^{N} O(U^k) \quad P[U] \propto e^{-S_E[U, \psi, \bar{\psi}]}$

A robust tool: lattice QCD

Non-perturbative 1st principles approach to Quantum Chromo Dynamics



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- Successful at T>0: static QCD properties⁴
 - (Pseudo)critical temperature: 154±9 MeV WB JHEP 1009 (2010) 073 - HotQCD PRD85 (2012) 054503
 - Equation of state as input for hydro-dynamics
 - Trace anomaly T^{μμ} = ε-3p strong coupling at T_C HotQCD PRD90 (2014) 094503 - WB PLB730 (2014) 99-104

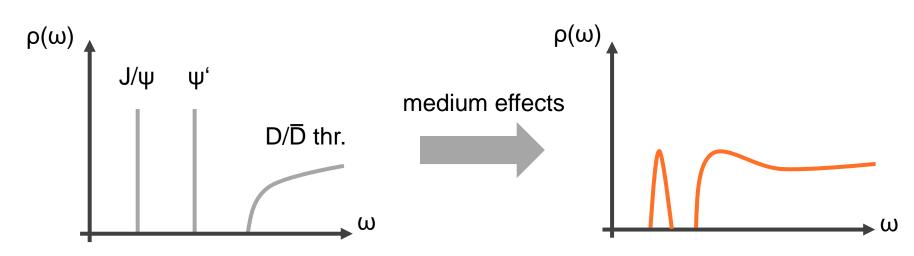


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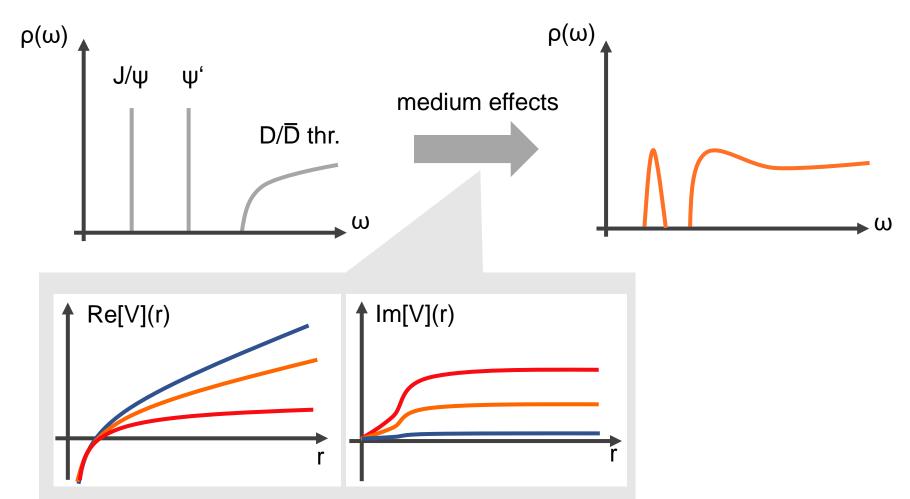
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Quarkonium in thermal equilibrium
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Intuition on in-medium modification from potential based studies



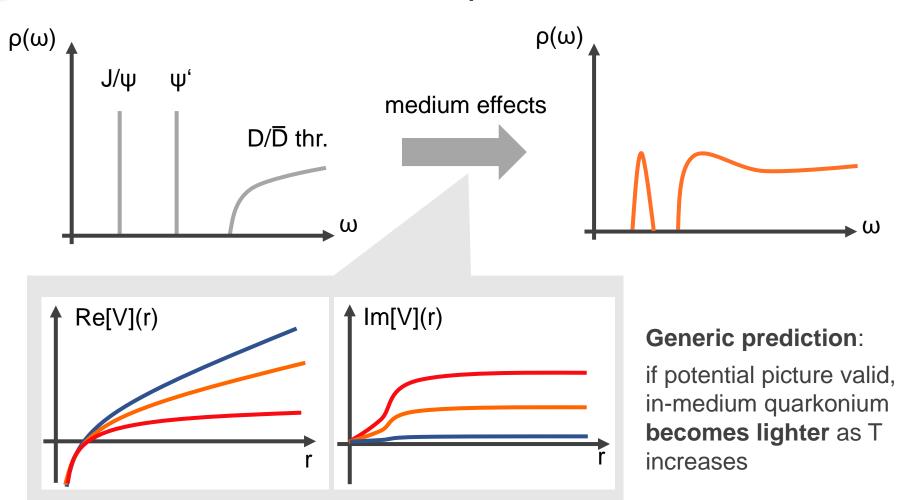
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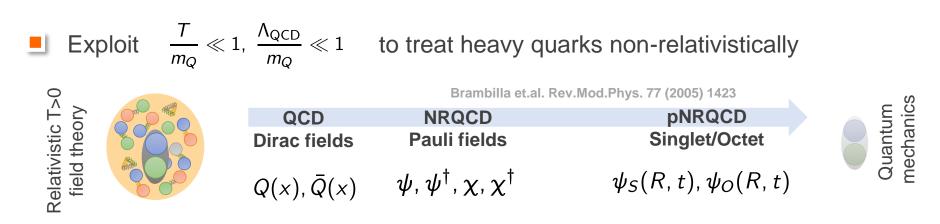
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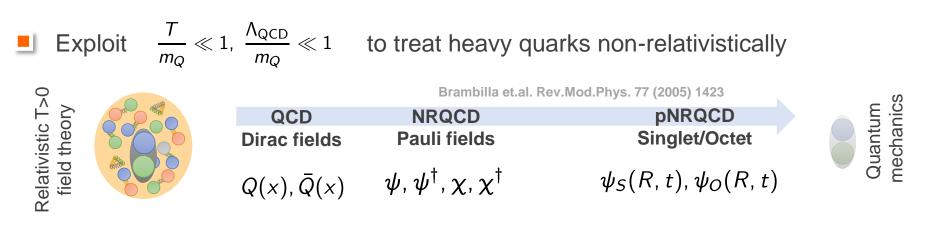
The QCD real-time interquark potential



$$i\partial_t \langle \psi_s(t)\psi_s(0)\rangle = \Big(V^{\rm QCD}(R) + \mathcal{O}(m_Q^{-1}) + \Theta(R,t)\Big) \langle \psi_s(t)\psi_s(0)\rangle$$

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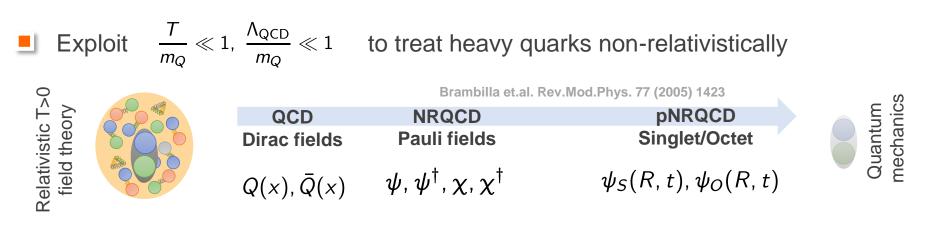
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Matching to underlying QCD in the infinite mass limit: Wilson loop

$$\langle \psi_{S}(R,t)\psi_{S}^{*}(R,0)\rangle_{pNRQCD} \equiv W_{\Box}(R,t) = \left\langle \operatorname{Tr}\left[\exp\left(-ig\int_{\Box}dx^{\mu}A_{\mu}(x)\right)\right]\right\rangle_{QCD}$$

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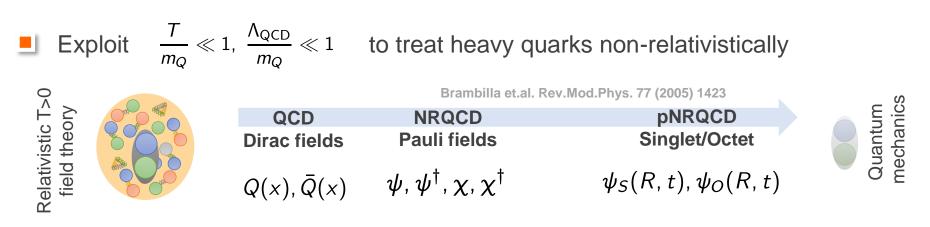
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Wilson loop: potential emerges at late times

$$V(R) = \lim_{t \to \infty} rac{i \partial_t W_{\Box}(R, t)}{W_{\Box}(R, t)}$$

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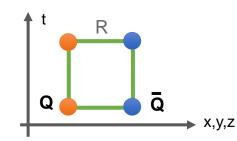
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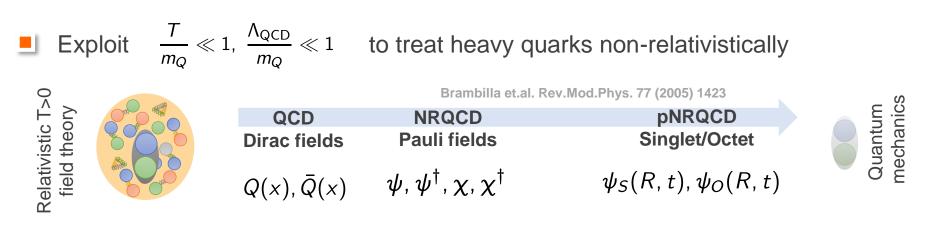
$$V(R) = \lim_{t \to \infty} \frac{i \partial_t W_{\Box}(R, t)}{W_{\Box}(R, t)} \in \mathbb{C}$$

Im[V]: Laine et al. JHEP03 (2007) 054; Beraudo et. al. NPA 806:312,2008 Brambilla et.al. PRD 78 (2008) 014017



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The QCD real-time interquark potential



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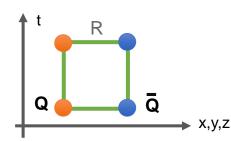
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In this form: Minkowski time quantities and not directly accessible on the lattice

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Inversion of Laplace transform required – highly ill-posed

How to connect to the Euclidean domain: **spectral functions**

$$W_{\Box}(R,t) = \int_{-\infty}^{\infty} d\omega \, e^{-i\omega t} \, \rho_{\Box}(R,\omega) \quad \longleftrightarrow \quad W_{\Box}(R,\tau) = \int_{-\infty}^{\infty} d\omega \, e^{-\omega \tau} \, \rho_{\Box}(R,\omega)$$

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Regularize this task using prior information – Bayes introduces prior P[p|I]=exp[S] M. Jarrell, J. Gubernatis, Physics Reports 269 (3) (1996)

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ho=
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Non-perturbative evaluation of V(R) University of Stavanger

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BR prior enforces: ρ positive definite, smoothness of ρ, result independent of units

Y.Burnier, A.R.
PRL 111 (2013) 18, 182003
$$S_{BR} = \alpha \int d\omega \left(1 - \frac{\rho}{m} + \log\left[\frac{\rho}{m}\right]\right)$$

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A.R., T.Hatsuda & S.Sasaki

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BR prior: better accuracy in sharp peak structures than MEM or BG but prone to ringing

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C.Fischer, J. Pawlowski,
A.R., C. Welzbacher
PRD98 (2018) 014009
$$S_{BR}^{smooth} = \alpha \int d\omega \left(\kappa \left(\frac{\partial \rho}{\partial \omega}\right)^2 + 1 - \frac{\rho}{m} + \log\left[\frac{\rho}{m}\right]\right)$$

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Latest results on the lattice potential



Lattices with dynamical u,d,s quarks (HISQ action, HotQCD & TUMQCD)

A. Bazavov et.al. PRD97 (2018) 014510, HotQCD PRD90 (2014) 094503

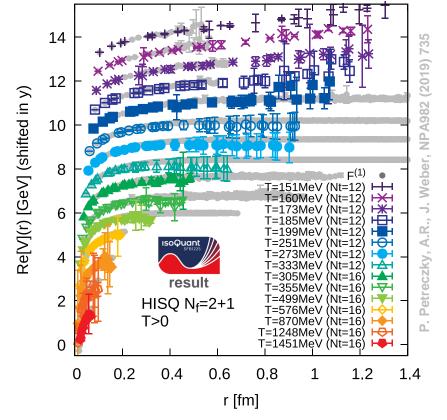
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- fixed box (N_s=48 N_T=12, N_T=16) & very high statistics 4000-9000 realizations
- Pade based extraction for Re[V] possible

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Smooth transition from Cornell @ T=0 to Debye screened @ T>T_c

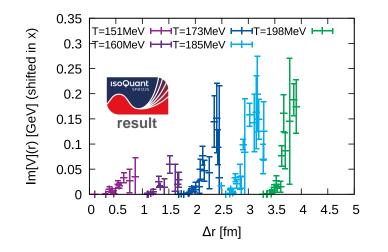


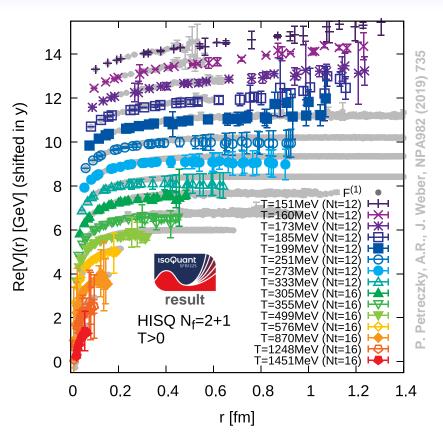
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Lattices with dynamical u,d,s quarks (HISQ action, HotQCD & TUMQCD)

A. Bazavov et.al. PRD97 (2018) 014510, HotQCD PRD90 (2014) 094503

- I realistic m_{π} ~161MeV (T=151-1451MeV)
- fixed box (N_s=48 N_T=12, N_T=16) & very high statistics 4000-9000 realizations
- Pade based extraction for Re[V] possible





- Smooth transition from Cornell @ T=0 to Debye screened @ T>T_c
- Finite Im[V] above T_c present



An improved Gauss law approach

For use in phenomenology applications: analytic expression for Re[V] and Im[V]

$$V_{Q\bar{Q}}^{T=0}(R) = V_C(R) + V_S(R) = -\frac{\alpha_S}{r} + \sigma r + c$$

Strategy:

 α_s, σ and c are vacuum prop. and do not change with T

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 $V(R) = aqR^a$

Coulombic: a=-1 q= α_s

$$ec
abla \Big(ec
abla V_{\mathcal{C}}(R)\Big) = -4\pi lpha_S \delta(ec R)$$

String-like: a=+1 q=
$$\sigma$$

 $\vec{\nabla}\left(\frac{\vec{\nabla}V_{S}(R)}{R^{2}}\right) = -4\pi\sigma\delta(\vec{R})$

V. V. Dixit, Mod. Phys. Lett. A 5, 227 (1990)

 C^{\prime}

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Immerse non-perturbative charge in weak coupling HTL medium: permittivity ε original idea: Y.Burnier, A.R. Phys.Lett. B753 (2016) 232 improved derivation D.Lafferty and A.R. arXiv:1906.00035

$$V^{med}(\mathbf{p}) = V^{vac}(\mathbf{p})/\epsilon(\mathbf{p}) \qquad \epsilon^{-1}(\vec{p}, m_D) = rac{p^2}{p^2 + m_D^2} - i\pi T rac{pm_D^2}{(p^2 + m_D^2)^2}$$

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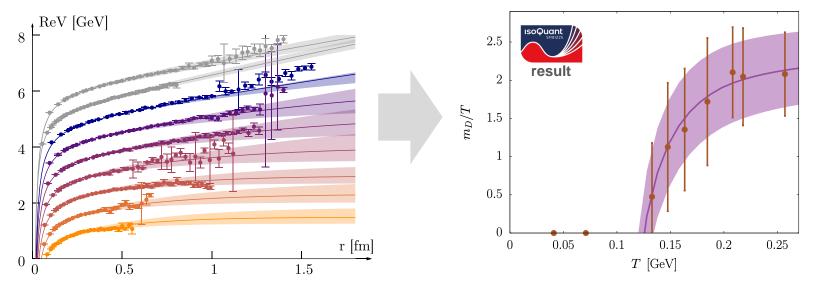
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3 vacuum parameters and 1 temperature dependent m_D fix both Re[V] and Im[V].

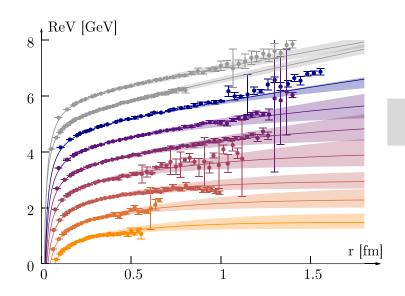


Gauss-Law result allows to fit Re[V] data even in the non-perturbative regime

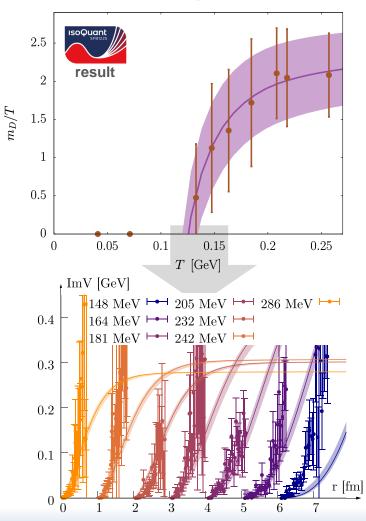


D.Lafferty and A.R. arXiv:1906.00035

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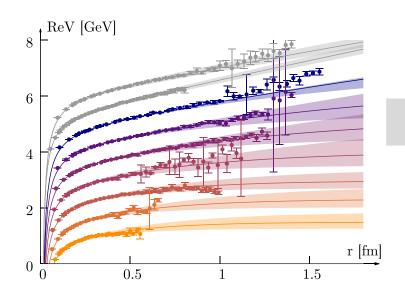


m_D defined from Re[V] allows to compute Gauss law prediction for Im[V]

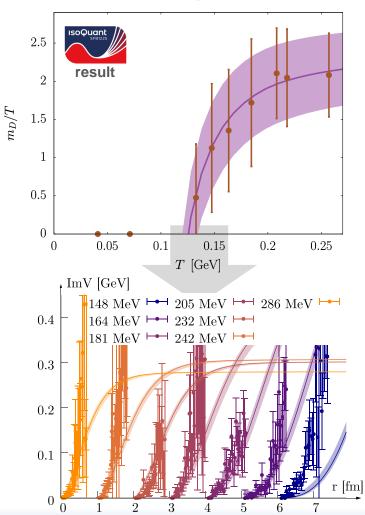


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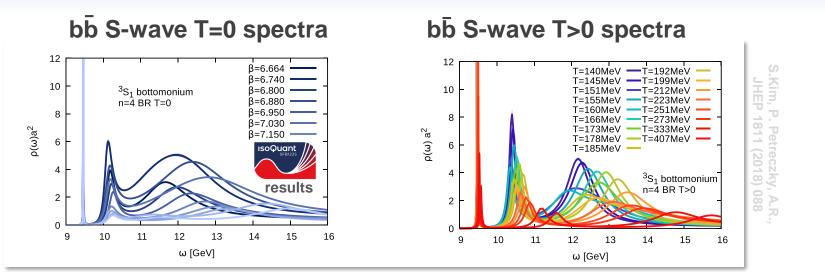


- Imp defined from Re[V] allows to compute Gauss law prediction for Im[V]
- recently extend the Gauss law to model quarkonium at finite velocity & µ_B



D.Lafferty and A.R. arXiv:1906.00035

Equilibrium spectral properties

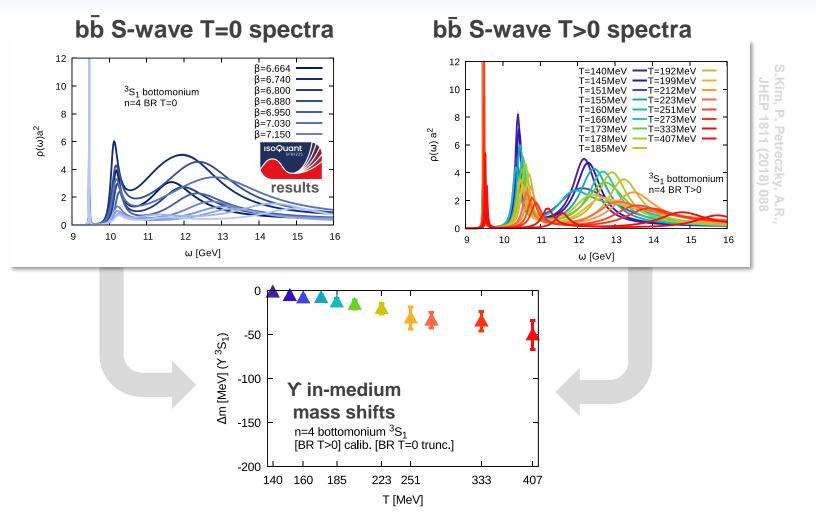


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Crucial step: **defining correct T=0 baseline** in presence of methods artifacts

Equilibrium spectral properties



University

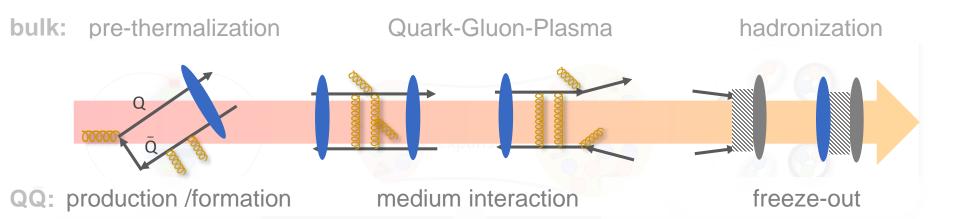
of Stavanger

Crucial step: **defining correct T=0 baseline** in presence of methods artifacts

For the first time consistent negative in medium mass shifts – ordered by E_{bind}

Open theory questions

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QQ realtime evol. in the initial stages

First exploratory steps in the glasma

classical statistical simulations for gluons & real-time NRQCD

(see poster by Alexander Lehmann)

Real-time QQ evol. in local thermal equilibrium

Beyond Schrödinger: **Open-quantum-systems** descr. of real-time evolution

Connecting OQS to EFT language of potential

Properties of equilibrium QQ

First principles extraction of the heavy quark potential

Novel phenomenological definition of the **Debye mass**

Extraction of thermal spectral properties on the lattice

with P. Petreczky, J. Weber: NPA982 (2019) 735 S. Kim, P. Petreczky, A.R. JHEP 1811 (2018) 088 with D. Lafferty arXiv:1906.00035

(with T. Miura, Y. Akamatsu, M. Asakawa arXiv:1908.06293)

ALEXANDER ROTHKOPF - UIS

The 1st QSEC Conference – September 24th 2019 – Heidelberg University – Germany

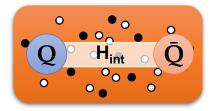
The open quantum systems picture

Need a general approach to describe quarkonium coupled to a thermal medium

Overall system is closed, hermitean Hamiltonian: von Neumann equation

$$H = H_{Qar{Q}} \otimes I_{med} + I_{Qar{Q}} \otimes H_{med} + H_{int}$$

$$\frac{d\rho}{dt} = -i[H,\rho]$$



Towards in-medium heavy quarkonium dynamics from first principles

The open quantum systems picture III University of Stavanger

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$$H_{\rm int} = \sum_m \Sigma_m \otimes \Xi_m$$

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The open quantum systems picture

Need a general approach to describe quarkonium coupled to a thermal medium

Solution Interest in the second secon

$$H = H_{Q\bar{Q}} \otimes I_{med} + I_{Q\bar{Q}} \otimes H_{med} + H_{int}$$
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Dynamics of the reduced QQbar system:

N. Brambilla et.al. arXiv:1903.08063 & PRD97 (2018) 074009, J.P. Blaizot, M. Escobedo JHEP 1806 (2018) 034

$$\rho_{Q\bar{Q}}=\mathsf{Tr}_{\mathit{med}}\big[\rho\big]$$

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$$\begin{array}{c} & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ \end{array} \quad \mathbf{Tr}_{\mathrm{med}} \quad \mathbf{V}(r) + ? \quad \mathbf{Q} \quad H_{\mathrm{int}} = \sum_{m} \Sigma_{m} \otimes \Xi_{m}$$

1

Dynamics of the reduced QQbar system: N. Brambilla et.al. arXiv:1903.08063 & PRD97 (2018) 074009,

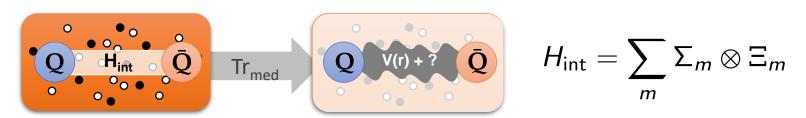
J.P. Blaizot, M. Escobedo JHEP 1806 (2018) 034

$$\rho_{Q\bar{Q}} = \operatorname{Tr}_{med}[\rho] \qquad \frac{d}{dt}\rho_{Q\bar{Q}} = \mathcal{V}\rho_{Q\bar{Q}}$$

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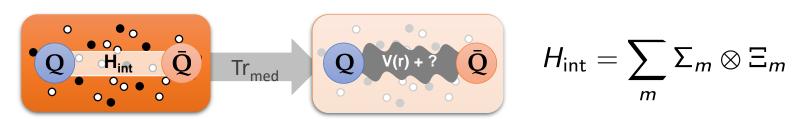
Separation of time-scales determines the nature of the e.o.m. :

 $\begin{array}{ll} \text{Environment relaxation scale } \tau_E : & \mbox{Q\bar{Q} system scale } \tau_S : & \mbox{Q\bar{Q} relaxation scale } \tau_{\text{rel}} : \\ & \langle \Xi_m(t) \Xi_m(0) \rangle \sim e^{-t/\tau_E} & \tau_S \sim 1/|\omega - \omega'| & \langle p(t) \rangle \propto e^{-t/\tau_{\text{rel}}} \end{array}$

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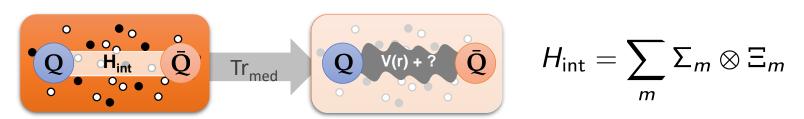
In case of Markovian time evolution ($au_E \ll au_{rel}$) leads to a Lindblad equation

$$\frac{d}{dt}\rho_{Q\bar{Q}} = -i[\tilde{H}_{Q\bar{Q}}, \rho_{Q\bar{Q}}] + \sum_{k} \gamma_{k} \left(L_{k}\rho_{Q\bar{Q}}L_{k}^{\dagger} - \frac{1}{2}L_{k}^{\dagger}L_{k}\rho_{Q\bar{Q}} - \frac{1}{2}\rho_{Q\bar{Q}}L_{k}^{\dagger}L_{k} \right)$$

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Feynman Vernon influence functional



Derivation via path integral formalism: **Feynman-Vernon influence functional** for details see Y. Akamatsu, Phys.Rev. D87 (2013) 4, 045016 and arXiv:1403.5783

$$\rho(t, x, y, X, Y) = \int dx_0 dy_0 dX_0 dY_0 \rho(0, x_0, y_0, X_0, Y_0) \int_{x_0, y_0, X_0, Y_0}^{x, y, X, Y} \mathcal{D}[\bar{x}, \bar{y}, \bar{X}, \bar{Y}] e^{iS[\bar{x}, \bar{X}] - iS[\bar{y}, \bar{Y}]}$$

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$$\rho_{Q\bar{Q}}(t,x,y) = \int dX dY \delta(X-Y) \rho(t,x,y,X,Y)$$

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I Approximate the FV functional to second order in the coupling & $au_{
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 $S_{FV} \approx S_{pot} [Re[V]] + S_{fluct} [Im[V]] + S_{diss} [Im[V]]$

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$$L_{\mathbf{k},a} = \sqrt{\frac{D(\mathbf{k})}{2}}\left[1 - \frac{\mathbf{k}}{4m_QT}\cdot\left(\frac{1}{2}\mathbf{P}_{CM} + \mathbf{p}\right)\right]e^{i\mathbf{k}\cdot\mathbf{r}/2}(T^a\otimes 1) - \sqrt{\frac{D(\mathbf{k})}{2}}\left[1 - \frac{\mathbf{k}}{4m_QT}\cdot\left(\frac{1}{2}\mathbf{P}_{CM} - \mathbf{p}\right)\right]e^{-i\mathbf{k}\cdot\mathbf{r}/2}(1\otimes T^a)$$

Feynman Vernon influence functional

Derivation via path integral formalism: **Feynman-Vernon influence functional** for details see Y. Akamatsu, Phys.Rev. D87 (2013) 4, 045016 and arXiv:1403.5783 $\rho(t, x, y, X, Y) = \int dx_0 dy_0 dX_0 dY_0 \rho(0, x_0, y_0, X_0, Y_0) \int_{x_0, y_0, X_0, Y_0}^{x, y, X, Y} \mathcal{D}[\bar{x}, \bar{y}, \bar{X}, \bar{Y}] e^{iS[\bar{x}, \bar{X}] - iS[\bar{y}, \bar{Y}]}$ $\rho_{Q\bar{Q}}(t, x, y) = \int dx_0 dy_0 \rho_{Q\bar{Q}}(0, x, y) \int_{x_0, y_0}^{x, y} \mathcal{D}[\bar{x}, \bar{y}] e^{iS_{Q\bar{Q}}[\bar{x}] - iS_{Q\bar{Q}}[\bar{y}] + \frac{iS_{FV}[\bar{x}, \bar{y}]}{medium - QQ interaction}}$

I Approximate the FV functional to second order in the coupling & $au_{
m E} \ll au_{
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$$L_{k,a} = \sqrt{\frac{D(k)}{2}}\left[1 - \frac{k}{4m_{Q}T}\cdot\left(\frac{1}{2}\mathbf{P}_{CM}+\mathbf{p}\right)\right]e^{i\mathbf{k}\cdot\mathbf{r}/2}(T^{a}\otimes1) - \sqrt{\frac{D(k)}{2}}\left[1 - \frac{k}{4m_{Q}T}\cdot\left(\frac{1}{2}\mathbf{P}_{CM}-\mathbf{p}\right)\right]e^{-i\mathbf{k}\cdot\mathbf{r}/2}(1\otimes T^{a})$$

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In QM language: Markovian evolution via quarkonium Lindblad equation

$$\frac{d}{dt}\rho_{Q\bar{Q}}(t) = -i\left[H_{Q\bar{Q}},\rho_{Q\bar{Q}}\right] + \sum_{i=1}^{N_{LB}}\gamma_{i}\left(\hat{L}_{i}\rho_{Q\bar{Q}}\hat{L}_{i}^{\dagger} - \frac{1}{2}\hat{L}_{i}\hat{L}_{i}^{\dagger}\rho_{Q\bar{Q}} - \frac{1}{2}\rho_{Q\bar{Q}}\hat{L}_{i}\hat{L}_{i}^{\dagger}\right) \qquad \stackrel{\tilde{D}(k) = g^{2}T\frac{\pi m_{D}^{2}}{k(k^{2} + m_{D}^{2})^{2}}$$
$$L_{k,a} = \sqrt{\frac{D(k)}{2}}\left[1 - \frac{k}{4m_{Q}T}\cdot\left(\frac{1}{2}\mathbf{P}_{CM}+\mathbf{p}\right)\right]e^{i\mathbf{k}\cdot\mathbf{r}/2}(T^{a}\otimes1) - \sqrt{\frac{D(k)}{2}}\left[1 - \frac{k}{4m_{Q}T}\cdot\left(\frac{1}{2}\mathbf{P}_{CM}-\mathbf{p}\right)\right]e^{-i\mathbf{k}\cdot\mathbf{r}/2}(1\otimes T^{a})$$

Full Lindblad dynamics cannot be described by deterministic Schrödinger equation

The stochastic potential



Based on scale separation & weak coupling approximation:

$$S_{FV} \approx S_{pot} [Re[V]] + S_{fluct} [Im[V]] + S_{diss} [Im[V]]$$

In QM language corresponds to Markovian evolution by Lindblad equation

$$\frac{d}{dt}\rho_{Q\bar{Q}}(t) = -i\left[H_{Q\bar{Q}},\rho_{Q\bar{Q}}\right] + \sum_{i=1}^{N_{LB}}\gamma_i\left(\hat{L}_i\rho_{Q\bar{Q}}\hat{L}_i^{\dagger} - \frac{1}{2}\hat{L}_i\hat{L}_i^{\dagger}\rho_{Q\bar{Q}} - \frac{1}{2}\rho_{Q\bar{Q}}\hat{L}_i\hat{L}_i^{\dagger}\right)$$

The stochastic potential



Based on scale separation & weak coupling approximation:

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Include the effects of fluctuations by taking the first order gradient expansion

The stochastic potential



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Include the effects of fluctuations by taking the first order gradient expansion

$$U = \exp\left[-i\int dt \left(-\frac{\nabla^2}{m_Q} + V(r) + \eta(t, \frac{\mathbf{r}}{2}) - \eta(t, -\frac{\mathbf{r}}{2})\right)\right]$$

$$S_{\text{pot}}$$

$$S_{\text{fluct}}$$

$$\langle \eta(t, \mathbf{x})\eta(t', \mathbf{y}) \rangle = \delta(t - t')D(\mathbf{x} - \mathbf{y})$$

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Fully unitary microscopic dynamics, no thermalization

The stochastic potential



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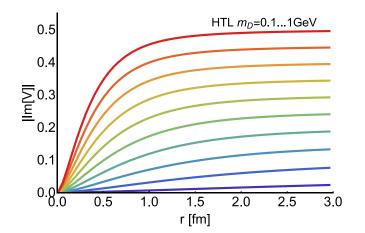
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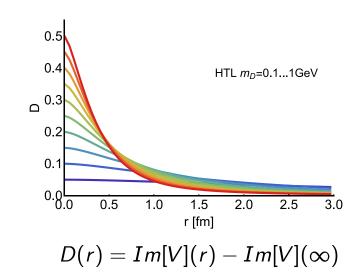
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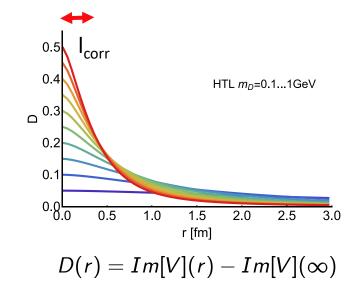
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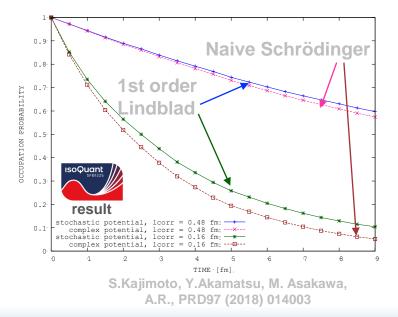
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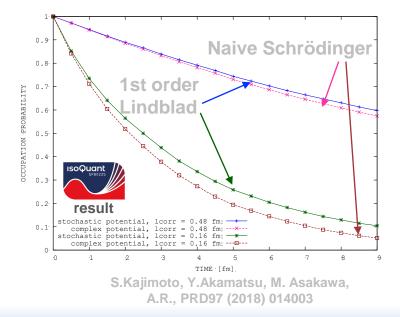
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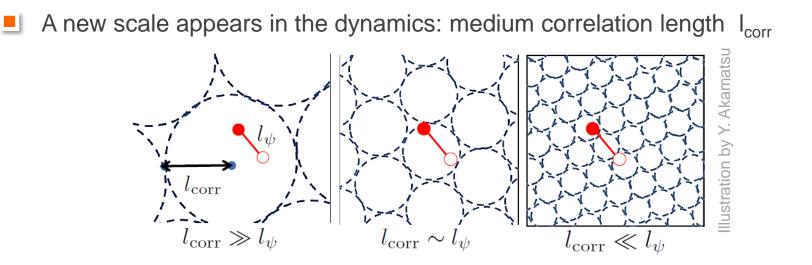
I Fully unitary microscopic dynamics, no thermalization $i\partial_t \langle \psi_{Q\bar{Q}}(t) \rangle = \left(-\frac{\nabla^2}{M} + Re[V] - i|Im[V]| \right) \langle \psi_{Q\bar{Q}}(t) \rangle$



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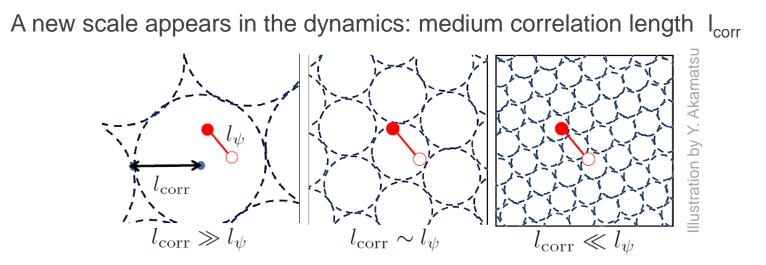
Decoherence



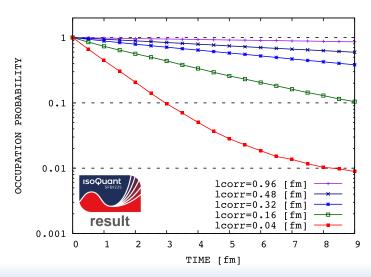


Decoherence





Fluctuations induce decoherence: select preferred basis – decay of populations





stability of quarkonium decreases: from Debye screening and fluctuations

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Towards full Lindblad dynamics (I)

University of Stavanger

Based on scale separation & weak coupling approximation:

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Unravel dynamics in wavefunction stochastic dynamics: Quantum State Diffusion T. Miura, Y.Akamatsu, M. Asakawa, A.R., arXiv:1908.06293

$$\begin{split} |d\psi\rangle &= |\psi(t+dt)\rangle - |\psi(t)\rangle \\ &= -iH|\psi(t)\rangle dt + \sum_{n} \begin{pmatrix} 2\langle L_{n}^{\dagger}\rangle\psi L_{n} - L_{n}^{\dagger}L_{n} \\ -\langle L_{n}^{\dagger}\rangle\psi \langle L_{n}\rangle\psi \end{pmatrix} |\psi(t)\rangle dt \\ &+ \sum_{n} \left(L_{n} - \langle L_{n}\rangle\psi\right) |\psi(t)\rangle d\xi_{n}, \end{split} \qquad L_{\mathbf{k},a} &= \sqrt{\frac{D(\mathbf{k})}{2}} \left[1 - \frac{\mathbf{k}}{4m_{Q}T} \cdot \left(\frac{1}{2}\mathbf{P}_{\mathsf{CM}} - \mathbf{p}\right)\right] e^{-i\mathbf{k}\cdot\mathbf{r}/2} (T^{a}\otimes 1) \\ &- \sqrt{\frac{D(\mathbf{k})}{2}} \left[1 - \frac{\mathbf{k}}{4m_{Q}T} \cdot \left(\frac{1}{2}\mathbf{P}_{\mathsf{CM}} - \mathbf{p}\right)\right] e^{-i\mathbf{k}\cdot\mathbf{r}/2} (1\otimes T^{a}) \\ &= \sum_{n} \left(L_{n} - \langle L_{n}\rangle\psi\right) |\psi(t)\rangle d\xi_{n}, \end{aligned} \qquad \tilde{D}(k) &= g^{2}T \frac{\pi m_{D}^{2}}{k(k^{2} + m_{D}^{2})^{2}}, \quad m_{D} = gT\sqrt{\frac{N_{c}}{3} + \frac{N_{f}}{6}} \\ &\text{weak coupling parameters, V(r) Debye form} \end{split}$$

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Towards full Lindblad dynamics (I)

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First "derivation" of phenomenological models based on nonlinear Schrödinger equation c.f. e.g. R. Katz, P. Gossiaux Annals Phys. 368 (2016) 267

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First genuine Lindblad implementation: previous works could not maintain positivity of ρ D. De Boni, JHEP 1708 (2017) 064

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Towards in-medium heavy quarkonium dynamics from first principles

Towards full Lindblad dynamics (II)

Based on scale separation & weak coupling approximation:

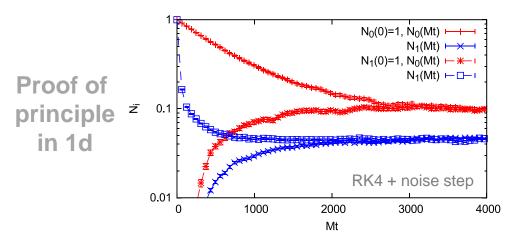
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Encouraging: admixtures become independent of initial conditions at late times

ALEXANDER ROTHKOPF - UIS The 1st QSEC Conditions– September 24th 2019 – Heidelberg University – Germany Towards in-medium heavy quarkonium dynamics from first principles

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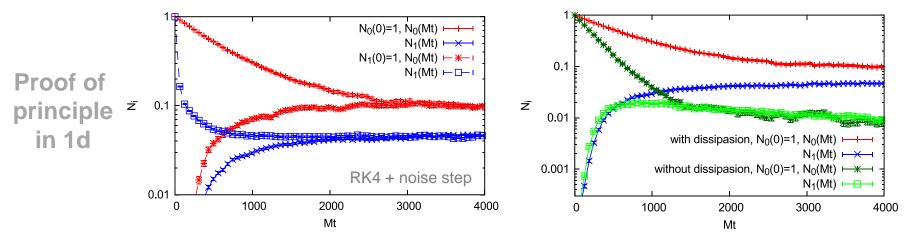
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Dissipative effects stabilize the ground state as they counteract separation of the pair **ALEXANDER ROTHKOPF - UIS** The 1st QSEC Conditions– September 24th 2019 – Heidelberg University – Germany



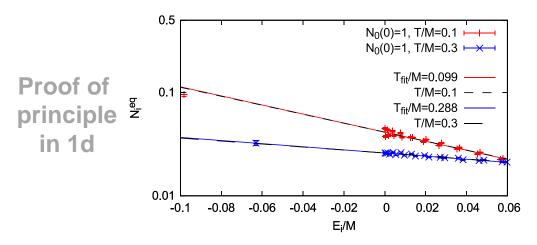
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Unravel dynamics in wavefunction stochastic dynamics: Quantum State Diffusion T. Miura, Y. Akamatsu, M. Asakawa,, A.R., arXiv:1908.06293



Distribution of states at late times agrees well with Boltzmann and yields consistent T

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Towards IN-MEDIUM HEAVY QUARKONIUM DYNAMICS FROM FIRST PRINCIPLES **Towards full Lindblad dynamics (III)**University of Stavanger

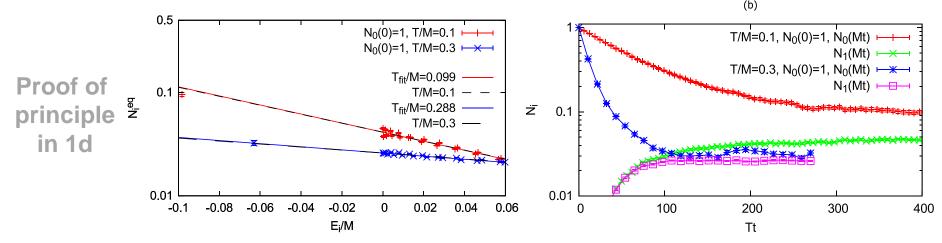
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Smaller m_Q leads to more efficient equilibration (decoherence more effective)

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Conclusion



- **Significant progress** in in-medium quarkonium theory
- Recent and ongoing studies on quarkonium dynamical properties
 - Novel extraction of the in-medium heavy quark potential on realistic lattices P. Petreczky, A.R., J. Weber, NPA982 (2019) 735
 - Improved analytic parametrization of V(R) using the generalized Gauss law D. Lafferty, A.R., arXiv:1906.00035
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with A. Lehmann (in preparation)

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Thank you for your attention

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